

Schwarzschild-de Sitter metric type solutions of a Nonlocal de Sitter Gravity Model

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Nonlocal Modification of GR

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Nonlocal square root gravity model

$$S = \frac{1}{16\pi G} \int_M \sqrt{R - 2\Lambda} F(\square) \sqrt{R - 2\Lambda} \sqrt{-g} d^4x,$$

where $F(\square) = 1 + \mathcal{F}(\square) = 1 + \sum_{n=1}^{\infty} f_n \square^n + \sum_{n=1}^{\infty} f_{-n} \square^{-n}$.

■ Construction

$$\begin{aligned} R - 2\Lambda &= \sqrt{R - 2\Lambda} \sqrt{R - 2\Lambda} \rightarrow \sqrt{R - 2\Lambda} F(\square) \sqrt{R - 2\Lambda} \\ &= R - 2\Lambda + \sqrt{R - 2\Lambda} \mathcal{F}(\square) \sqrt{R - 2\Lambda} \end{aligned}$$

Equations of motion

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Equation of motion are

$$-\frac{1}{2}g_{\mu\nu}\sqrt{R-2\Lambda}\mathcal{F}(\square)\sqrt{R-2\Lambda} + R_{\mu\nu}W - K_{\mu\nu}W + \frac{1}{2}\Omega_{\mu\nu} = -(G_{\mu\nu} + \Lambda g_{\mu\nu})$$

$$\Omega_{\mu\nu} = \sum_{n=1}^{\infty} f_n \sum_{l=0}^{n-1} S_{\mu\nu}(\square^l \sqrt{R-2\Lambda}, \square^{n-1-l} \sqrt{R-2\Lambda})$$
$$- \sum_{n=1}^{\infty} f_{-n} \sum_{l=0}^{n-1} S_{\mu\nu}(\square^{-(l+1)} \sqrt{R-2\Lambda}, \square^{-(n-l)} \sqrt{R-2\Lambda}),$$

$$K_{\mu\nu} = \nabla_{\mu} \nabla_{\nu} - g_{\mu\nu} \square,$$

$$S_{\mu\nu}(A, B) = g_{\mu\nu} \nabla^{\alpha} A \nabla_{\alpha} B - 2 \nabla_{\mu} A \nabla_{\nu} B + g_{\mu\nu} A \square B,$$

$$W = \frac{1}{\sqrt{R-2\Lambda}} \mathcal{F}(\square) \sqrt{R-2\Lambda}.$$

Eigenvalue problem

If we assume $\square\sqrt{R-2\Lambda} = p\sqrt{R-2\Lambda}$, EOM are simplified to

$$W = \mathcal{F}(p)$$

$$\Omega_{\mu\nu} = \mathcal{F}'(p)S_{\mu\nu}(\sqrt{R-2\Lambda}, \sqrt{R-2\Lambda}),$$

$$(G_{\mu\nu} + \Lambda g_{\mu\nu})(1 + \mathcal{F}(q)) + \frac{1}{2}\mathcal{F}'(q)S_{\mu\nu}(\sqrt{R-2\Lambda}, \sqrt{R-2\Lambda}) = 0.$$

It is evident that EOM are satisfied if $\mathcal{F}(p) = -1$ and $\mathcal{F}'(p) = 0$.

Schwarzschild-de Sitter-type metric

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We want to investigate our model outside the spherically symmetric massive body. Since this model is a nonlocal generalization of general relativity with the cosmological constant Λ , it is natural to consider a generalization of the Schwarzschild-de Sitter metric starting from the standard Schwarzschild expression

$$ds^2 = -A(r)dt^2 + \frac{1}{A(r)}dr^2 + r^2d\theta^2 + r^2 \sin^2 \theta d\varphi^2.$$

The corresponding scalar curvature R of above metric is

$$R = \frac{2 - 2A(r) - 4rA'(r) - r^2A''(r)}{r^2} = \frac{1}{r^2} \frac{\partial^2}{\partial r^2} [r^2(1 - A(r))].$$

As we concluded earlier, to find a solution of EoM it is sufficient to solve an eigenvalue problem $\square\sqrt{R - 2\Lambda} = q\sqrt{R - 2\Lambda}$. Note that here d'Alembertian \square acts in the following way:

$$\square u(r) = A(r) u''(r) + (A'(r) + \frac{2}{r} A(r)) u'(r) = \frac{1}{r^2} \frac{\partial}{\partial r} \left[r^2 A(r) \frac{\partial u}{\partial r} \right],$$

where $u(r)$ is any differentiable scalar function.

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where $u(r)$ is any differentiable scalar function. Let us now consider function $A(r)$ in the form

$$A(r) = 1 - \frac{\mu}{r} - \frac{\nu}{r^2} - \frac{\Lambda}{6} r^2 - f(r),$$

where μ and ν are some parameters to be discussed later. Then one can show that for our choice of $A(r)$ holds

$$R(r) = \frac{1}{r^2} \frac{\partial^2}{\partial r^2} [r^2(1 - A(r))] = 2\Lambda + \frac{1}{r^2} \frac{\partial^2}{\partial r^2} [r^2 f(r)].$$

Denoting $u(r) = \sqrt{R - 2\Lambda}$ and using previous results we obtain

$$\square \sqrt{R - 2\Lambda} = \frac{1}{r^2} \frac{\partial}{\partial r} \left[r^2 A(r) \frac{\partial}{\partial r} \sqrt{R - 2\Lambda} \right] = q \sqrt{R - 2\Lambda}.$$

To get an approximative solution we take $A(r) \approx 1$, what is applicable when

$$\left| \frac{\mu}{r} \right| \ll 1, \quad \left| \frac{\nu}{r^2} \right| \ll 1, \quad |\Lambda r^2| \ll 1, \quad |f(r)| \ll 1,$$

Under these conditions \square operator becomes a Laplacian Δ

$$\Delta\sqrt{R-2\Lambda} = \frac{1}{r^2} \frac{\partial}{\partial r} \left[r^2 \frac{\partial}{\partial r} \sqrt{R-2\Lambda} \right] = q \sqrt{R-2\Lambda}.$$

One can easily find eigenfunctions of Δ operator

$$u(r) = \frac{C_1}{r} e^{\sqrt{q} r} + \frac{C_2}{r} e^{-\sqrt{q} r}.$$

The next step is looking for general solution of

$$R(r) - 2\Lambda = \frac{1}{r^2} \frac{\partial^2}{\partial r^2} [r^2 f(r)] = u^2(r),$$

which is

$$f(r) = -\frac{\Lambda r^2}{6} - \frac{C_1^2}{4qr^2} e^{2\sqrt{q}r} - \frac{C_2^2}{4qr^2} e^{-2\sqrt{q}r} - C_1 C_2 + \frac{C_3}{r} + \frac{C_4}{r^2}.$$

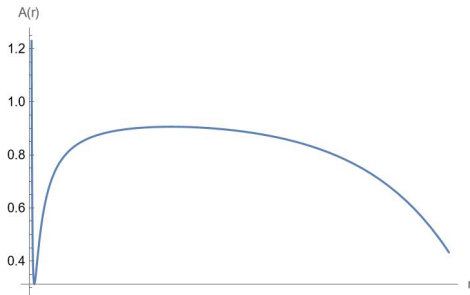
Then the value of $A(r)$ and $R(r)$ become

$$A(r) = 1 - \frac{\mu}{r} - \frac{\nu}{r^2} - \frac{\Lambda r^2}{6} - \frac{C_1^2}{4qr^2} e^{2\sqrt{q}r} - \frac{C_2^2}{4qr^2} e^{-2\sqrt{q}r} - C_1 C_2,$$

$$R(r) = 2\Lambda + \left(\frac{C_1}{r} e^{\sqrt{q}r} + \frac{C_2}{r} e^{-\sqrt{q}r} \right)^2.$$

$$A(r) = 1 - \frac{\mu}{r} - \frac{\nu}{r^2} - \frac{\Lambda r^2}{6} - \frac{C_1^2}{4qr^2} e^{2\sqrt{q}r} - \frac{C_2^2}{4qr^2} e^{-2\sqrt{q}r} - C_1 C_2$$

where q , C_1 and C_2 are free parameters and their values should be determined by observations.



Some relevant references

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