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### Predictions in Cartan F(R) Gravity

T. I., M. Taniguchi, Symmetry 14, 1830 (2022), T. I., H. Sakamoto, M. Taniguchi, arXiv:2304.14769 [gr-qc] to be appear in JCAP.



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#### Outline

- Modified theory of gravity
- Cartan F(R) gravity
- Spacetime evolution in Cartan F(R) gravity
- Fluctuations in CMB
- Conclusion

# Modified theory of gravity

Einstein's General Relativity

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$



Why?

Einstein's General Relativity

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$



Non-linear equation

Einstein's General Relativity

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

 Phenomenological consequences Apsidal precession of the planet Mercury, Dense stars, Black holes, Gravitational lens, Gravitational wave, Expansion of the universe,

Einstein's General Relativity

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Phenomenological consequences
 Apsidal precession of the planet Mercury,
 Dense stars, Black holes, ←Heavy neutron stars
 Gravitational lens, ←Dark matter
 Gravitational wave, ←Just started
 Expansion of the universe, ←Accelerated expansion
 ··· ←Small scale < 0.01mm</p>

#### Heavy neutron stars



http://xtreme.as.arizona.edu/NeutronStars/

#### Dark matter and dark energy



$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Action

$$S = \frac{1}{2\kappa} \int d^4x \sqrt{-g} (R - 2\Lambda) + S_{\text{matter}}$$



$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Action

Modified Gravity

- Higher order terms
- Non-local terms
- Gauss-Bonnet

• Torsion

$$S = \frac{1}{2\kappa} \int d^4x \sqrt{-g(R)} - 2\Lambda) + S_{\text{matter}}$$

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Particle physics

- Neutrino
- Axion
- Super partners

• • • •

Action

$$S = \frac{1}{2\kappa} \int d^4x \sqrt{-g} (R - 2\Lambda) + S_{\text{matter}}$$

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Spacetime

• Extra dimensions

. . .

Action

$$S = \frac{1}{2\kappa} \int \sqrt[4]{g} \sqrt{-g} (R - 2\Lambda) + S_{\text{matter}}$$

M. Montesinos, R. Romero and D. Gonzalez, Class. Quant. Grav. 37 (2020) 045008.
T.P. Sotiriou and S. Liberati, J. Phys. Conf. Ser. 68 (2007) 012022.
T.P. Sotiriou and S. Liberati, Annals Phys. 322 (2007) 935.
D. Iosifidis, A.C. Petkou and C.G. Tsagas, Gen. Rel. Grav. 51 (2019) 66.
S. Capozziello, R. Cianci, C. Stornaiolo and S. Vignolo, Class. Quant. Grav. 24 (2007) 6417.
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T.P. Sotiriou, Class. Quant. Grav. 26 (2009) 152001.
S. Capozziello and S. Vignolo, Annalen Phys. 19 (2010) 238.
G.J. Olmo, Int. J. Mod. Phys. D 20 (2011) 413.

# Cartan F(R) gravity

T. I., M. Taniguchi, Symmetry 14, 1830 (2022).

#### Modified gravity in Cartan formalism

• Action

$$\begin{split} S &= \frac{1}{2\kappa} \int d^4x \sqrt{-g} (R-2\Lambda) + S_{\rm matter} \qquad g = \det g_{\mu\nu} \\ & & \\ S &= \frac{1}{2\kappa} \int d^4x e F(R) + S_{\rm matter} \qquad e = \det e^i{}_{\mu} \end{split}$$

- Einstein-Cartan geometry
- Modified Lagrangian density

#### Cartan formalism E. Cartan (1923), T. W. B. Kibble (1961), D. W. Sciama(1962)

• Vierbein

Flat (local Lorentz frame)
Curved (general coordinate frame)

$$\underline{g_{\mu\nu}} = \underline{\eta_{ij}} e^i{}_{\mu} e^j{}_{\nu}$$

Covariant Derivative

$$\nabla_{\nu} e^{k}{}_{\mu} = \partial_{\nu} e^{k}{}_{\mu} + \omega^{k}{}_{l\nu} e^{l}{}_{\mu} - \Gamma^{\lambda}{}_{\mu\nu} e^{k}{}_{\lambda} = 0.$$
Spin connection Affine connection

• Torsion

$$T^{\lambda}{}_{\mu\nu} = \Gamma^{\lambda}{}_{\mu\nu} - \Gamma^{\lambda}{}_{\nu\mu}$$





#### Nontorsion and contorsion parts

• Affine connection

$$\Gamma^{\lambda}{}_{\mu\nu} = (\Gamma_{E})^{\lambda}{}_{\mu\nu} + K^{\lambda}{}_{\mu\nu}$$
Nontorsion Contorsion

$$(\Gamma_{E})^{\lambda}{}_{\mu\nu} \equiv \frac{1}{2}g^{\lambda\rho}(\partial_{\mu}g_{\nu\rho} + \partial_{\nu}g_{\rho\mu} - \partial_{\rho}g_{\mu\nu}$$

$$K^{\lambda}{}_{\mu\nu} = \frac{1}{2}\left(T^{\lambda}{}_{\mu\nu} + T_{\mu\nu}{}^{\lambda} + T_{\nu\mu}{}^{\lambda}\right)$$

#### Field equations

Variations of action

 $\delta S/\delta e^{i}{}_{\mu} = 0$   $\Rightarrow$  Modified Einstein equation Here, we assume that  $F'R^{i}{}_{\mu} - \frac{1}{2}e^{i}{}_{\mu}F(R) = M_{\mathrm{Pl}}{}^{-2}\Sigma^{i}{}_{\mu} \to R(\Sigma)$ the matter Lagrangian does not depend on the spinor connection.  $\delta S/\delta w^{i}{}_{j\mu} = 0 \implies$  Cartan equation  $T^{\mu}{}_{kl} - e_{l}{}^{\mu}T_{k} + e_{k}{}^{\mu}T_{l} + (e_{k}{}^{\alpha}e_{l}{}^{\mu} - e_{k}{}^{\mu}e_{l}{}^{\alpha})\partial_{\alpha}\ln F'(R) = 0$  $T^{k}{}_{ij} = \frac{1}{2} (\delta^{k}{}_{j} e_{i}{}^{\lambda} - \delta^{k}{}_{i} e_{j}{}^{\lambda}) \partial_{\lambda} \ln F'(R(\Sigma))$ 

#### Scalar Tensor theory

T. I., M. Taniguchi, Symmetry 14, 1830 (2022)

• We introduce a scalaron field

$$\phi \equiv -\sqrt{\frac{3}{2}}M_{\rm Pl}\ln F'(R)$$

and rewrite the action (without any conformal transformation)

$$S = \int d^4 x e \left( \frac{M_{\rm Pl}^2}{2} R_E - \frac{1}{2} \partial_\lambda \phi \partial^\lambda \phi - V(\phi) \right)$$
  

$$V(\phi) = - \left. \frac{M_{\rm Pl}^2}{2} (F(R) - R) \right|_{R=R(\phi)}$$
  
We drop a total derivative term.

Spacetime evolution in Cartan F(R) gravity

#### Expanding Universe



C. Faucher-Giguère, A. Lidz, and L. Hernquist, Science 319, 5859 (47)

#### Energy source for accelerated expansion

- Homogeneous and isotropic spacetime  $ds^2 = c^2 dt^2 + a^2(t)(dx^2 + dy^2 + dz^2)$
- Energy density



Radiation	$a(t) \propto t^{1/2}$
Matter	$a(t) \propto t^{2/3}$
Potential energy	$a(t) \propto over(at)$
Cosmological const.	$u(\iota) \propto \exp(\alpha \iota)$



#### Quasi de-Sitter expansion

• Friedmann eq.



A. D. Linde, Contemp. Concepts Phys. 5, 1 (1990).

#### Exit from inflation

• Equation of motion

$$\ddot{\varphi} + 3\frac{\dot{a}}{a}\dot{\varphi} = -\frac{\partial V}{\partial\varphi}$$

• Deceleration parameter

$$q \equiv -\frac{a\ddot{a}}{\dot{a}^2} \to 0$$



A. D. Linde, Contemp. Concepts Phys. 5, 1 (1990).

#### Slow-roll Inflation

• Slow-roll parameters

$$3\left(\frac{\dot{a}}{a}\right)^2 = \frac{1}{2}\dot{\phi}^2 + V(\phi), \ \dot{\phi}^2 \ll V(\phi)$$

 $V(\Phi)$  behaves as constant at  $\Lambda_{inf}$ 

$$\epsilon = \frac{M_{\rm Pl}^2}{2} \left(\frac{V_{\phi}}{V}\right)^2, \ \eta = \frac{M_{\rm Pl}^2 V_{\phi\phi}}{V}$$

 $\left\{ \begin{array}{l} \text{During inflation } \epsilon,\eta \ll 1 \\ \text{Exit from inflation } \epsilon,\eta \sim 1 \end{array} \right.$ 



# Fluctuations in CMB (cosmic microwave background)

T. I., H. Sakamoto, M. Taniguchi, arXiv:2304.14769 [gr-qc] to be appear in JCAP.



#### Quantum fluctuations

 $\varphi + \delta \varphi \\ \to \mathcal{P}_s(k)$ 

Scalar type fluctuation Origin: quantum fluctuation of scalar field

Tensor type fluctuation Origin: quantum fluctuation of space-time

 $g^{\mu\nu} + \delta h^{\mu\nu} \\ \to \mathcal{P}_t(k)$ 

#### Observed CMB fluctuations

• Rescaled scalar type fluctuation

$$\mathcal{P}_s(k) \equiv A_s \left(\frac{k}{k_0}\right)^{n_s - 1}$$

 Rescaled tensor type fluctuation

$$\mathcal{P}_t(k) \equiv A_t \left(\frac{k}{k_0}\right)^{n_t}$$

- Tensor to scalar ratio  $r \equiv \frac{\mathcal{P}_t(k)}{\mathcal{P}_s(k)}$ 

$$n_s - 1 = -6\epsilon + 2\eta$$
  
 $r = 16\epsilon$   
expressed by  $\epsilon, \eta$ 

#### Power law model

E. Elizalde, S. Nojiri, S. D. Odintsov, L. Sebastiani and S. Zerbini (2011)

Modified theory of gravity

$$F(R) = R - \frac{R^n}{(2nM^2)^{n-1}n}$$

Scalaron potential

$$V(\phi) = -\frac{M_{\rm Pl}^2}{2} f(R)|_{R=R(\phi)}$$
$$= M_{\rm Pl}^2 M^2 (1 - e^{-\sqrt{\frac{2}{3}}\frac{\phi}{M_{\rm Pl}}})^{1 + \frac{1}{n-1}}$$



#### CMB fluctuations ( $n_s$ and r)



#### Logarithmic model

S. Nojiri, S. D. Odintsov (2014)

- Modified theory of gravity  $F(R) = R \alpha R \ln \left( \frac{R}{R_0} + 1 \right)$
- Scalaron potential



#### CMB fluctuations ( $n_s$ and r)



#### CMB fluctuations ( $n_s$ and r)



T. I., H. Sakamoto, M. Taniguchi, arXiv:2304.14769 [gr-qc] to be appear in JCAP.

#### After the inflation era (logarithmic model)



#### After the inflation era (logarithmic model)



1+z

# Conclusion

#### Summary

- We study a modified theory of gravity on Einstein-Cartan geometry.
- The Cartan F(R) can be rewritten as a scalar tensor theory without conformal transformation. T. I., M. Taniguchi, Symmetry 14, 1830 (2022).
- The inflaton potential energy can induce the inflationary expansion of the universe.
- We observed a robustness of the CMB fluctuations.

T. I., H. Sakamoto, M. Taniguchi, arXiv:2304.14769 [gr-qc] to be appear in JCAP.

• There is a possibility to induce the current accelerated expansion.

#### Open questions

- Many models of Cartan F(R) gravity induce the inflationary expansion. How to distinguish the models?
- Matter fields may have an decisive role for the spacetime structure. What is a role of the matter fields?