

Multidimensional solitons

Boris A. Malomed

**Department of Physical Electronics, School of Electrical
Engineering, Faculty of Engineering
Tel Aviv University, Tel Aviv, Israel**

(1) Introduction

Formation of *localized structures* in the form of *solitons* in many physical settings is accounted for by the interplay of nonlinear *self-attraction* (alias *self-focusing*) of physical fields (e.g., electromagnetic waves in photonics, or macroscopic wave functions, alias *matter waves*, in *Bose-Einstein condensates*, **BECs**) and basic linear effects, such as *dispersion* or *diffraction*.

Arguably, the most important model which creates solitons is the *nonlinear Schrödinger* (NLS) *equation*:

$$i \frac{\partial \Psi}{\partial t} = -\frac{1}{2} \frac{\partial^2 \Psi}{\partial x^2} - |\Psi|^2 \Psi.$$

It generates a family of bright-soliton solutions with two free parameters - amplitude η and velocity c :

$$\Psi(x, t) = e^{icx + i(\eta^2 - c^2)t/2} \frac{\eta}{\cosh(\eta(x - ct))}.$$

The **exact integrability** of this equation, and of its counterpart with the **self-repulsive (defocusing)** nonlinearity,

$$i \frac{\partial \Psi}{\partial t} = -\frac{1}{2} \frac{\partial^2 \Psi}{\partial x^2} \pm |\Psi|^2 \Psi,$$

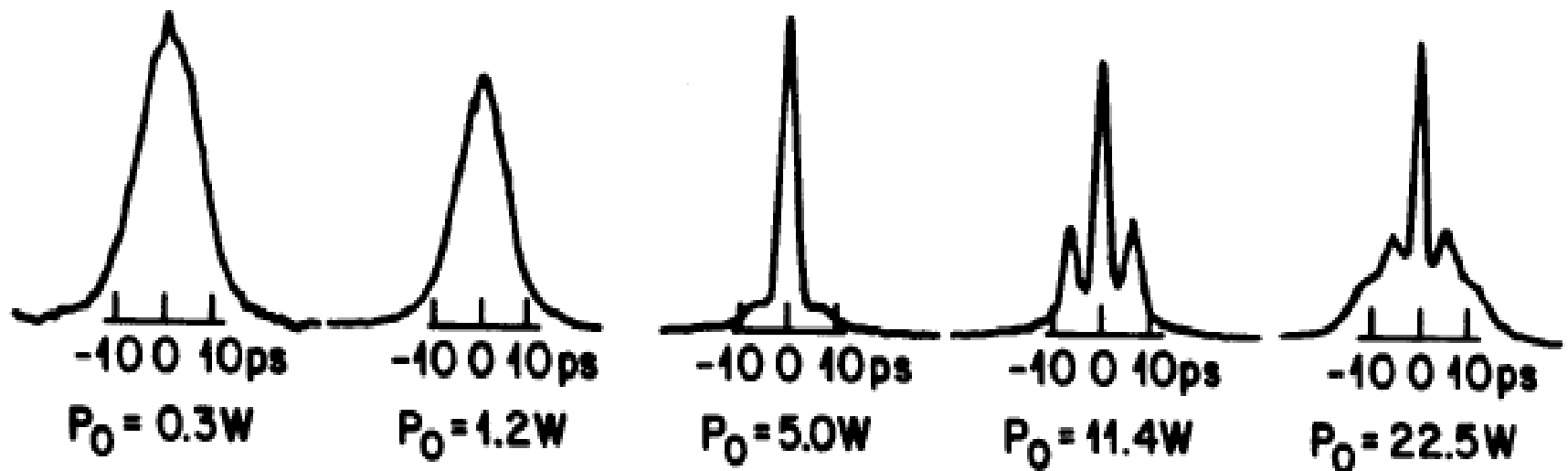
by means of the **inverse-scattering transform**, had been discovered in the classical work by **V.E. Zakharov** and **A.B. Shabat** (originally published in Russian),

V. E. Zakharov and A. B. Shabat, *Exact Theory of Two-dimensional Self-focusing and One-dimensional Self-modulation of Waves in Nonlinear Media*,

J. Exp. Theor. Phys. **34**, 62-69 (1972).

Solitons of the NLS type in *nonlinear optical fibers* were predicted by Hasegawa and Tappert in 1973 (Appl. Phys. Lett. **23**, 142-144), and experimentally created by Mollenauer, Stolen and Gordon in 1980 (Phys. Rev. Lett. **45**, 1095-1098).

They also observed two-soliton bound states (*breathers*):

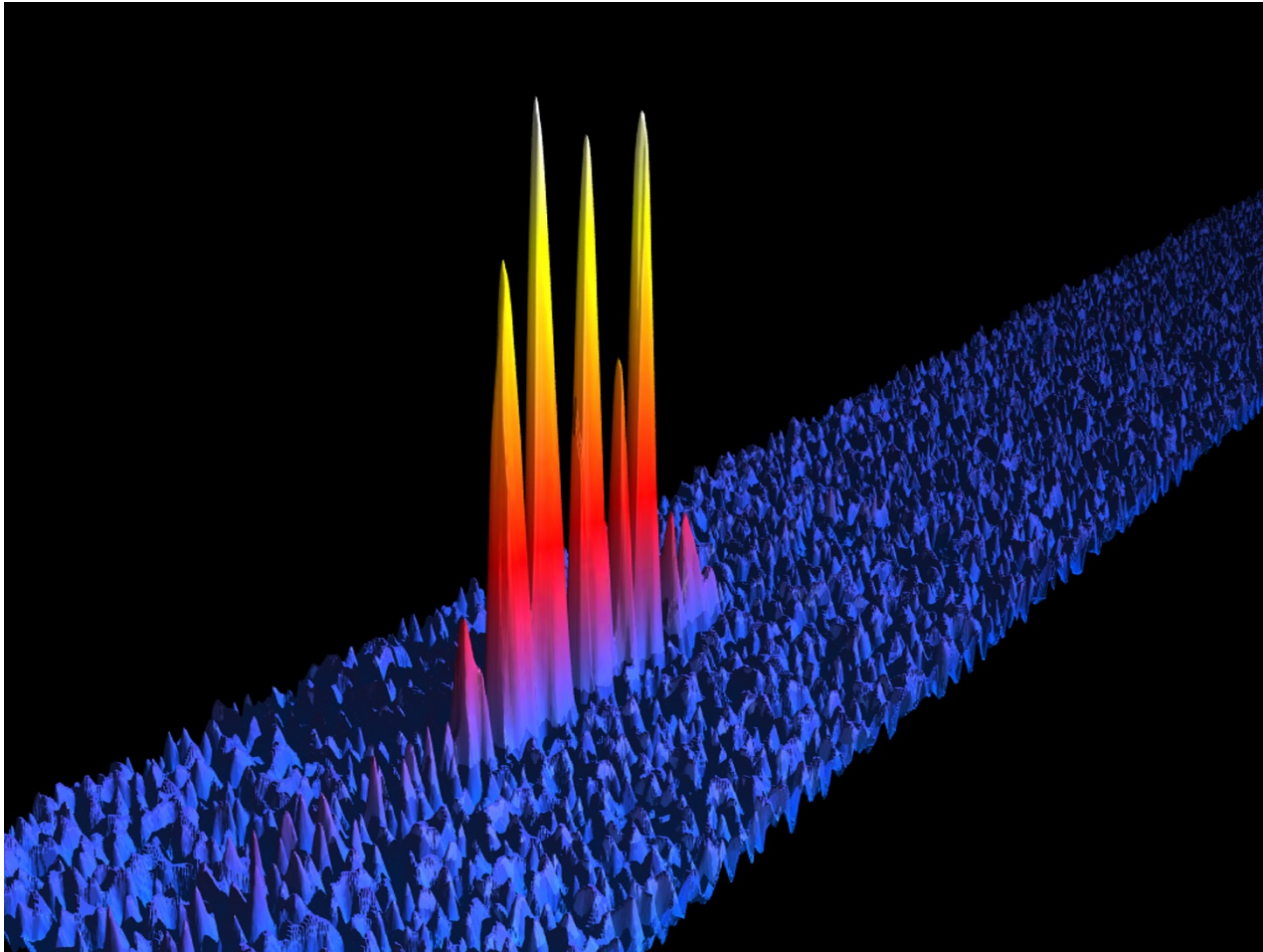


Another famous realization of solitons was demonstrated in **BEC** loaded in nearly one-dimensional (“cigar-shaped”) trapping potentials. Such *bright matter-wave solitons* were first created in the condensate of ^7Li atoms:

K.E. Strecker, G.B. Partridge, A.G. Truscott, and R. G. Hulet, Nature **417**, 150 (2002);

L. Khaykovich, F. Schreck, G. Ferrari, T. Bourdel, J. Cubizolles, L.D. Carr, Y. Castin, and C. Salomon, Science **296**, 1290 (2002).

The famous experimental picture of the atomic density distribution in a chain of **7 matter-wave solitons** with unequal amplitudes in ^7Li (from the work of R. Hulet *et al.*):



(2) The subject of the talk: multidimensional solitons

A more challenging objective is to **predict** and **create experimentally** **two-** and **three-dimensional** (**2D** and **3D**) solitons. On the contrary to their **1D** counterparts, which are, normally, ***stable***, the ubiquitous ***cubic self-focusing*** always makes **2D** and **3D** solitons ***unstable***, as ***exactly the same*** self-interaction gives rise to fundamental ***destabilizing effects***: the ***critical collapse*** in **2D**, and ***supercritical collapse*** in **3D** (i.e., ***catastrophic self-compression*** of the wave function, which leads to ***formation of a singularity*** after a finite evolution time).

First, let us look at **2D** solitons produced by the **NLS** equation with the **self-focusing cubic nonlinearity**,

$$iu_t + (1/2)(u_{xx} + u_{yy}) + |u|^2u = 0.$$

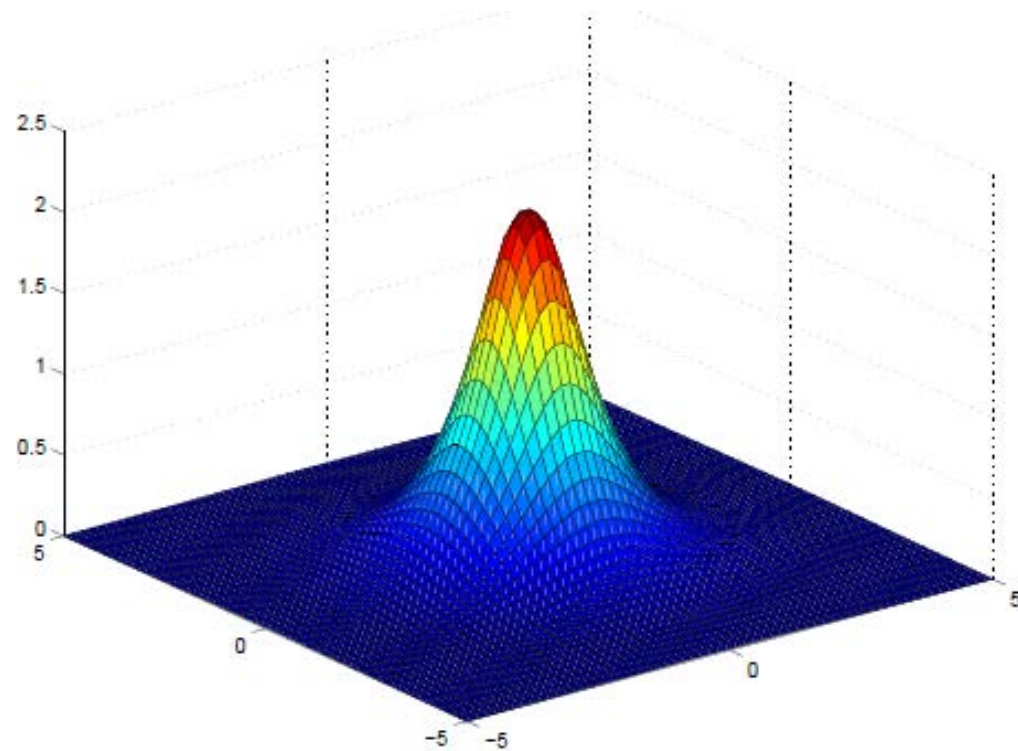
Soliton solutions with integer **vorticity** S and **chemical potential** μ are looked for, in the polar coordinates (r, θ) , as

$$u = \exp(-i\mu t + iS\theta) U_s(r).$$

The real radial wave function $U_s(r)$ satisfies the equation

$$\mu U_s + \frac{1}{2} \left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{S^2}{r^2} \right) U_s + U_s^3 = 0.$$

The profile of $U_0(r)$ for a **fundamental (zero-vorticity, $S = 0$) 2D soliton** [alias the *Townes' soliton*, R.Y. Chiao, E. Garmire & C. H. Townes, Phys. Rev. Lett . **13**, 479 (1964)]:



The families of the Townes' solitons (and their vortex counterparts) are **degenerate**: due to the specific **conformal symmetry** of the **2D NLS** equation with the cubic nonlinearity, the norm of each family (with **$S = 0, 1, 2$, etc.**),

$$N_S = 2\pi \int_0^\infty U_S^2(r; \mu) r dr,$$

takes a **single value**, which **does not** depend on the soliton's **chemical potential**, μ . In particular, for $S = 0$ (the **fundamental Townes' solitons**), $N_0 \approx 5.85$.

For $S \geq 1$, the norm is accurately approximated by

$$N_s = 2\pi \int_0^\infty U_s^2(r; \mu) r dr \approx 4\sqrt{3} \pi S,$$

as demonstrated in

PHYSICAL REVIEW A **94**, 053611 (2016)

Stable giant vortex annuli in microwave-coupled atomic condensates

Jieli Qin and Guangjiong Dong

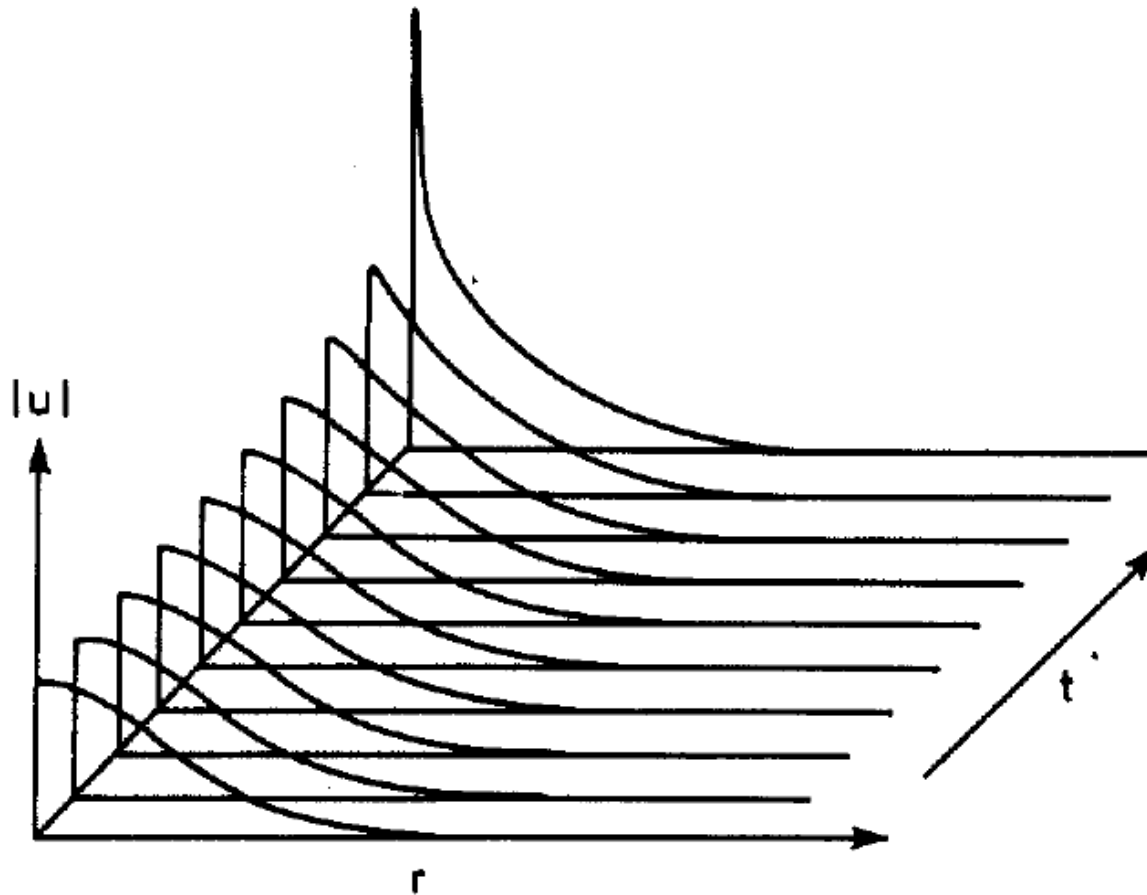
*State Key Laboratory of Precision Spectroscopy, Department of Physics, East China Normal University, Shanghai 200062, China
and Collaborative Innovation Center of Extreme Optics, Shanxi University, Taiyuan, Shanxi 030006, China*

Boris A. Malomed

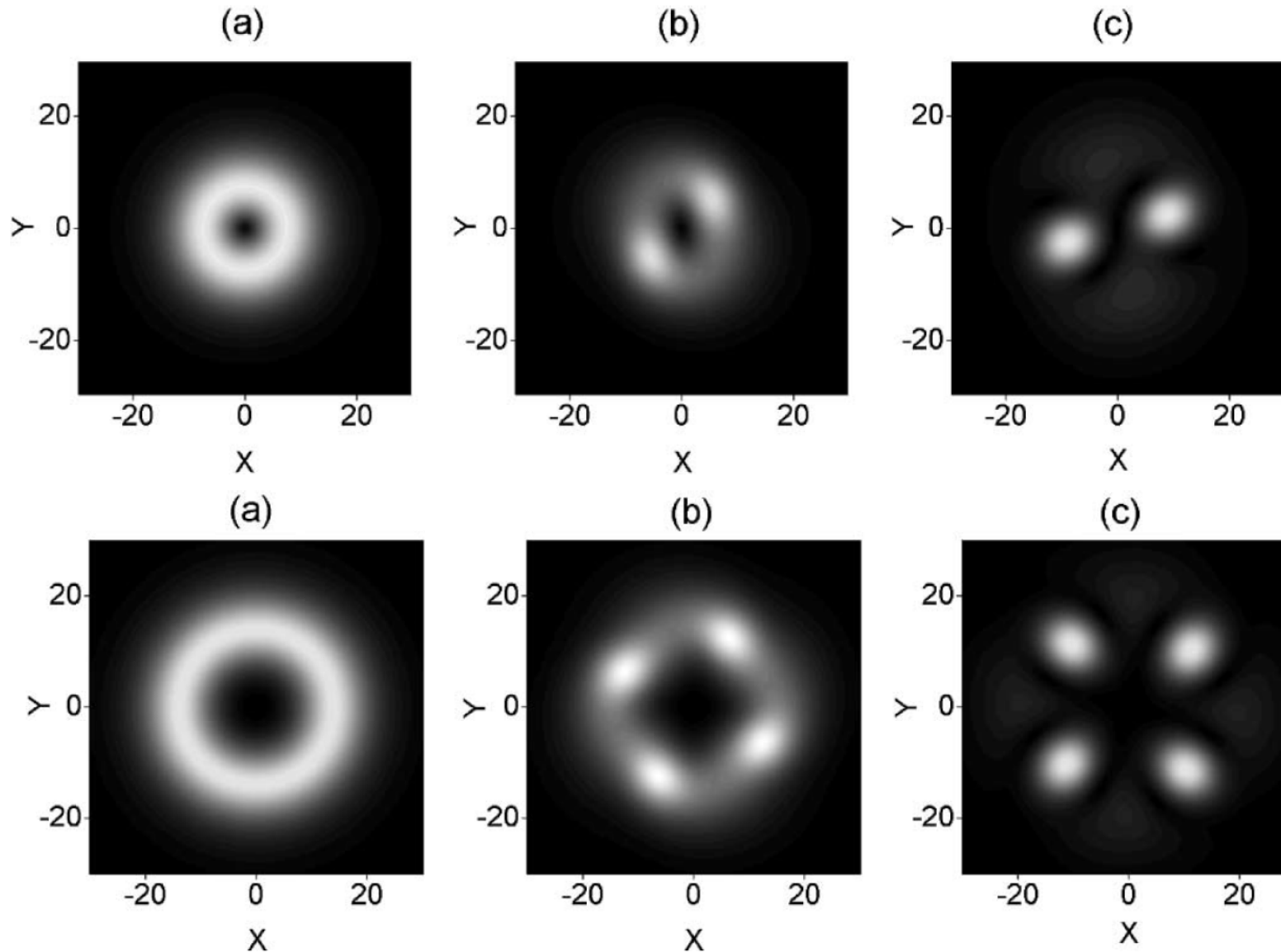
*Department of Physical Electronics, School of Electrical Engineering, Faculty of Engineering, Tel Aviv University, Ramat Aviv 69978, Israel
and Laboratory of Nonlinear-Optical Informatics, ITMO University, St. Petersburg 197101, Russia*

(Received 3 March 2016; revised manuscript received 17 October 2016; published 14 November 2016)

The numerically simulated development of the ***collapse*** of the fundamental Townes soliton in **2D** (the evolution of the radial cross section of the collapsing soliton towards the formation of a ***singularity***):



Examples of the *spontaneous splitting* of *unstable* 2D *vortex solitons* with vorticities $S = 1$ and $S = 2$:



The norm of the **Townes' soliton** is a **threshold value** which separates **collapsing** ($N > N_0$) and **decaying** ($N < N_0$) localized solutions of the **2D NLS** equation.

A **stabilizing mechanism**, added to the **2D NLS** equation, may act by letting the norm of **2D** solitons take values $N < N_0$ (but **avoiding the decay**), so that they **cannot** undergo the **collapse**, as their norm falls below the threshold value necessary to initiate the collapse.

Stabilization of **vortex solitons** is a **still more challenging problem**.

However, in the **3D NLS** equation, the situation is more difficult, as the collapse is **supercritical**, with **zero threshold**. This fact suggests that **3D** solitons may be made **metastable**, in the best case, as **the collapse remains possible**.

Thus, **stabilization** is the **most important** issue for **multidimensional solitons** (alias **spatiotemporal solitons**, in terms of optics). A review summarizing theoretical and experimental findings in this area:

Nature Reviews Physics 1, 185-197 (2019)

Frontiers in multidimensional self-trapping of nonlinear fields and matter

Yaroslav V. Kartashov^{1,2}, Gregory E. Astrakharchik³, Boris A. Malomed^{4,5,6,7}
and Lluís Torner^{1,8} *

Another review article, specifically focused on **2D** and **3D** solitons with *embedded vorticity* (which is one of central topics of the present talk):

Physica D 399 (2019) 108–137



Contents lists available at [ScienceDirect](#)

Physica D

journal homepage: www.elsevier.com/locate/physd

Invited Review Article

(INVITED) Vortex solitons: Old results and new perspectives

Boris A. Malomed

Department of Physical Electronics, School of Electrical Engineering, Faculty of Engineering, and Center for Light–Matter Interaction, Tel Aviv University, Tel Aviv 39040, Israel

A recent review article summarizes theoretical and experimental findings for multidimensional solitons in *dissipative media*:

Chaos, Solitons and Fractals 163 (2022) 112526



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Chaos, Solitons and Fractals

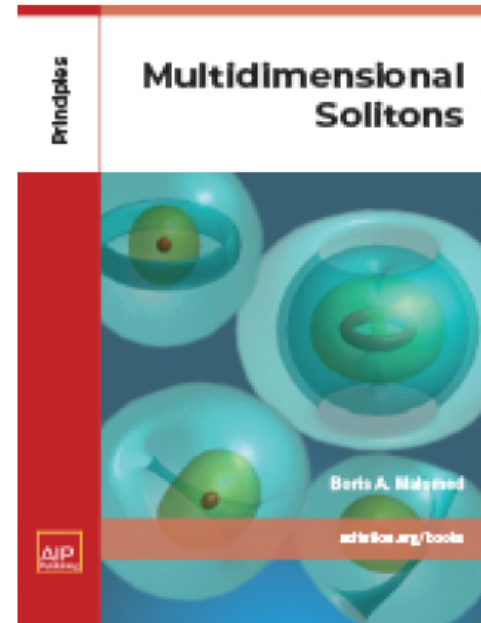
journal homepage: www.elsevier.com/locate/chaos


Multidimensional dissipative solitons and solitary vortices

B.A. Malomed^{*}

*Department of Physical Electronics, School of Electrical Engineering, Faculty of Engineering, Tel Aviv University, Tel Aviv 69978, Israel
Instituto de Alta Investigación, Universidad de Tarapacá, Casilla 7D, Arica, Chile*

All the results, theoretical and experimental ones, are presented in a systematic form in a **new book** (published by the American Institute of Physics, Melville, NY, 2022):



 Full . September 2022

Multidimensional Solitons

Author Boris A. Malomed

Publisher: AIP Publishing

Pages: 406

Copyright year: 2022

As the **stability** is the main issue for **2D** and **3D** settings, the rest of the talk is structured according to basic mechanisms which provide for **stabilization** of the multidimensional solitons:

- (3) Relatively old findings: **Stable 2D and 3D vortex solitons** in models with the **cubic-quintic (CQ)** nonlinearity (in a brief form).
- (4) A relatively new model: **Stable 2D and 3D composite solitons** in **two-component spin-orbit (SO)-coupled BEC**.
- (5) **Latest** theoretical and experimental results: the prediction and creation of **3D and 2D matter-wave solitons** (“*quantum droplets*”) stabilized by **quantum fluctuations**.
- (6) Conclusions.

(3) 2D and 3D Systems with the cubic-quintic nonlinearity

The stabilization of **2D** and **3D** fundamental and vortical solitons can be provided by a combination of ***competing self-focusing cubic*** and ***self-defocusing quintic*** nonlinear terms.

In optics, the **3D NLS** equation can be written as the equation governing the *spatiotemporal evolution* of the electromagnetic field:

$$i \frac{\partial u}{\partial z} + \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial \tau^2} + |u|^2 u - |u|^4 u = 0.$$

Stationary solutions with integer vorticity, $S \geq 0$, and propagation constant (wavenumber) k , are looked for in the *cylindrical spatiotemporal coordinates* as

$$u(x, y, \tau, z) = U(r, \tau) \exp(ikz + iS\theta), \quad r \equiv \sqrt{x^2 + y^2},$$

with $U(r, \tau)$ satisfying the stationary equation:

$$\left(\frac{\partial^2 U}{\partial r^2} + \frac{1}{r} \frac{\partial U}{\partial r} - \frac{S^2}{r^2} U + \frac{\partial^2 U}{\partial \tau^2} \right) + U^3 - U^5 = kR.$$

The 2D reduction of the **3D** equations implies dropping the term $\partial^2 U / \partial \tau^2$.

The stability of *fundamental solitons* ($\mathbf{S} = \mathbf{0}$) in the framework of this equation is obvious. A nontrivial problem is the *stability of vortex solitons* against splitting by azimuthal perturbations. For **2D vortex solitons**, the *stability* was first reported in

J. Opt. Soc. Am. B/Vol. 14, No. 8/August 1997

M. Quiroga-Teixeiro and H. Michinel

Stable azimuthal stationary state in quintic nonlinear optical media

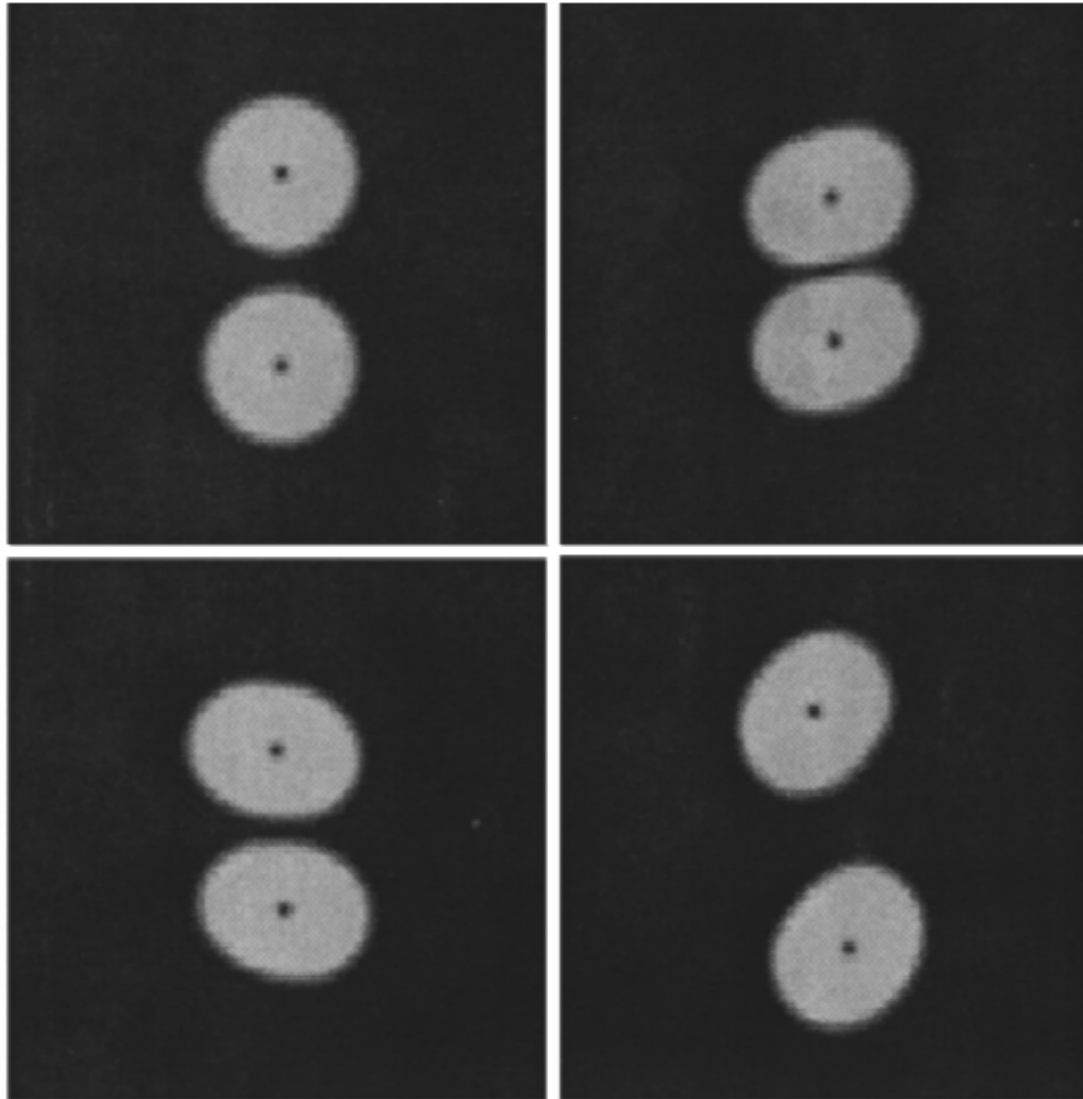
M. Quiroga-Teixeiro

Institute for Electromagnetic Field Theory, Chalmers University of Technology, S-412 96, Göteborg, Sweden

H. Michinel

Departamento de Física Aplicada, Escola Universitaria de Óptica, Universidade de Santiago de Compostela, E-157 06, Santiago de Compostela, Galicia, Spain

An example of the **dynamical stability** and **elastic collision** of **2D vortex solitons** with $S = 1$:



Experimentally, the stability of **(2+1)D** *fundamental* ($S = 0$) solitons in an optical *cubic-quintic* medium was demonstrated relatively recently:

PRL **110**, 013901 (2013)

PHYSICAL REVIEW LETTERS

week ending
4 JANUARY 2013

Robust Two-Dimensional Spatial Solitons in Liquid Carbon Disulfide

Edilson L. Falcão-Filho* and Cid B. de Araújo

Departamento de Física, Universidade Federal de Pernambuco, 50670-901 Recife, Pernambuco, Brazil

Georges Boudebs, Hervé Leblond, and Vladimir Skarka

LUNAM Université, Université d'Angers, Laboratoire de Photonique d'Angers, EA 4464, 49045 Angers, France

(Received 29 March 2012; published 2 January 2013)

The excitation of near-infrared $(2 + 1)$ D solitons in liquid carbon disulfide is demonstrated due to the simultaneous contribution of the third- and fifth-order susceptibilities. Solitons propagating free from diffraction for more than 10 Rayleigh lengths although damped, were observed to support the proposed soliton behavior. Numerical calculations using a nonlinear Schrödinger-type equation were also performed.

Fully stable *vortex solitons* have not yet been created in the experiment. However, *quasi-stable* ones have been reported in

PHYSICAL REVIEW A **93**, 013840 (2016)

Robust self-trapping of vortex beams in a saturable optical medium

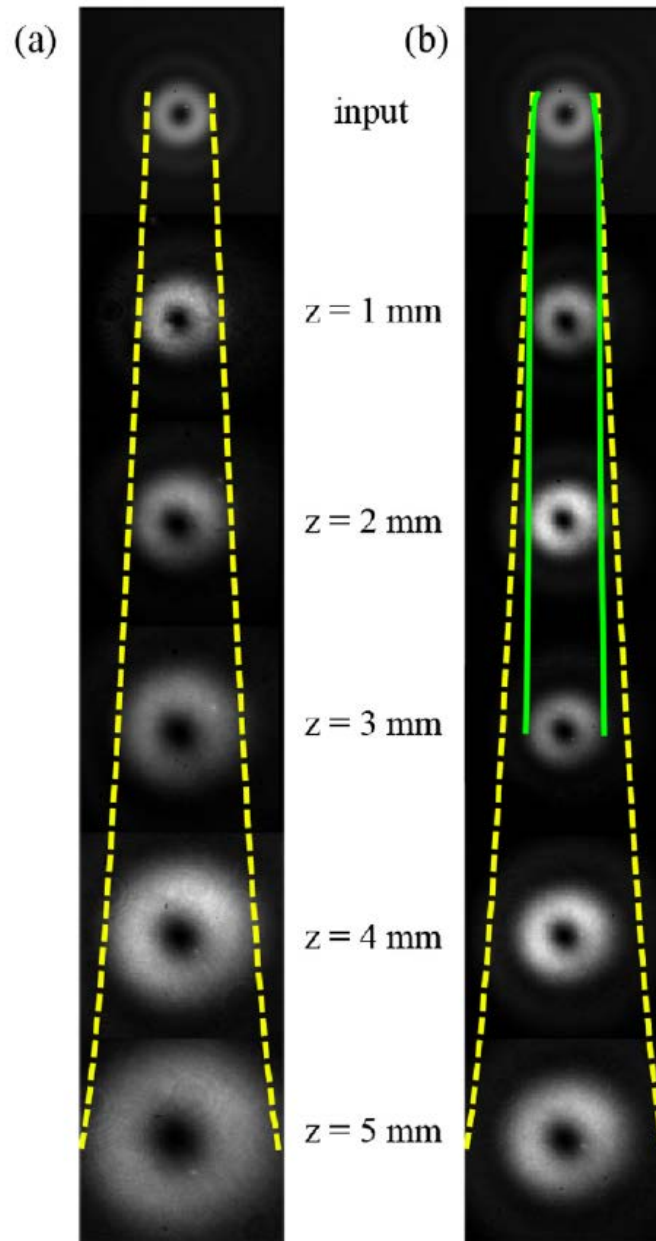
Albert S. Reyna,^{1,*} Georges Boudebs,² Boris A. Malomed,^{1,†} and Cid B. de Araújo¹

¹*Departamento de Física, Universidade Federal de Pernambuco, 50670-901, Recife, Pernambuco, Brazil*

²*LUNAM Université, Université d'Angers, Laboratoire de Photonique d'Angers, EA 4464, 49045 Angers, France*

(Received 28 March 2015; published 21 January 2016)

The vortex soliton starts expansion due to nonlinear losses (multiphoton absorption).



A challenging problem is to construct **3D** vortex solitons in the cubic-quintic medium, and analyze their **stability**. Theoretically, this was done long ago in:

Stable Spinning Optical Solitons in Three Dimensions

D. Mihalache,^{1,2,5} D. Mazilu,^{1,2} L.-C. Crasovan,^{1,5} I. Towers,^{3,4} A. V. Buryak,³ B. A. Malomed,⁴ L. Torner,⁵
J. P. Torres,⁵ and F. Lederer²

¹*Department of Theoretical Physics, Institute of Atomic Physics, P.O. Box MG-6, Bucharest, Romania*

²*Institute of Solid State Theory and Theoretical Optics, Friedrich-Schiller University Jena, Max-Wien-Platz 1,
D-07743, Jena, Germany*

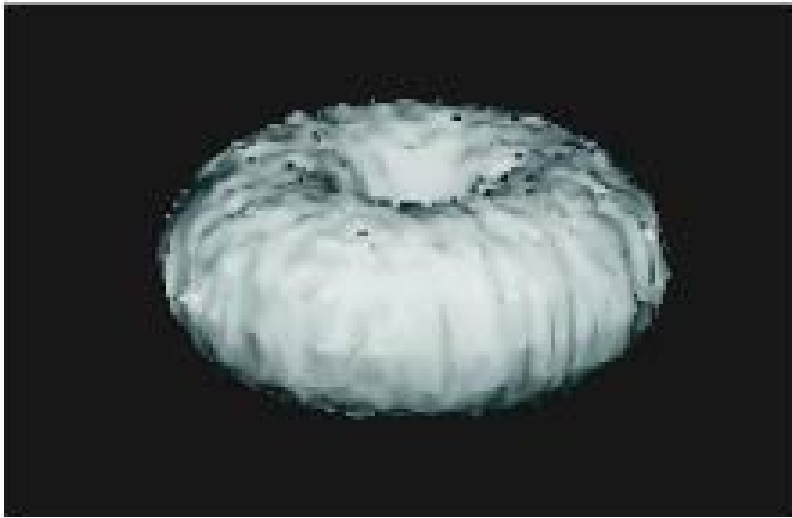
³*School of Mathematics and Statistics, University of New South Wales at ADFA, Canberra, ACT 2600, Australia*

⁴*Faculty of Engineering, Tel Aviv University, Tel Aviv 69978, Israel*

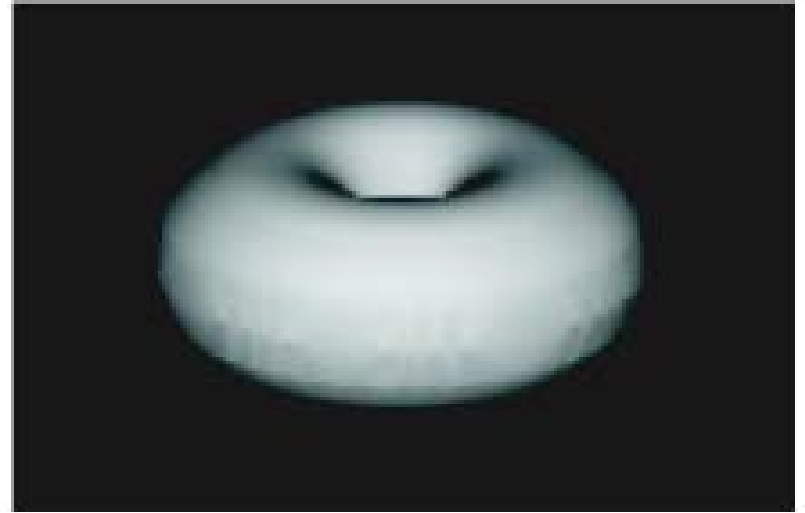
⁵*Department of Signal Theory and Communications, Universitat Politècnica de Catalunya, ES 08034 Barcelona, Spain*
(Received 2 July 2001; published 4 February 2002)

An example of numerically simulated *recovery* of a strongly perturbed **3D** *stable* soliton with intrinsic vorticity $S = 1$ (*doughnut*):

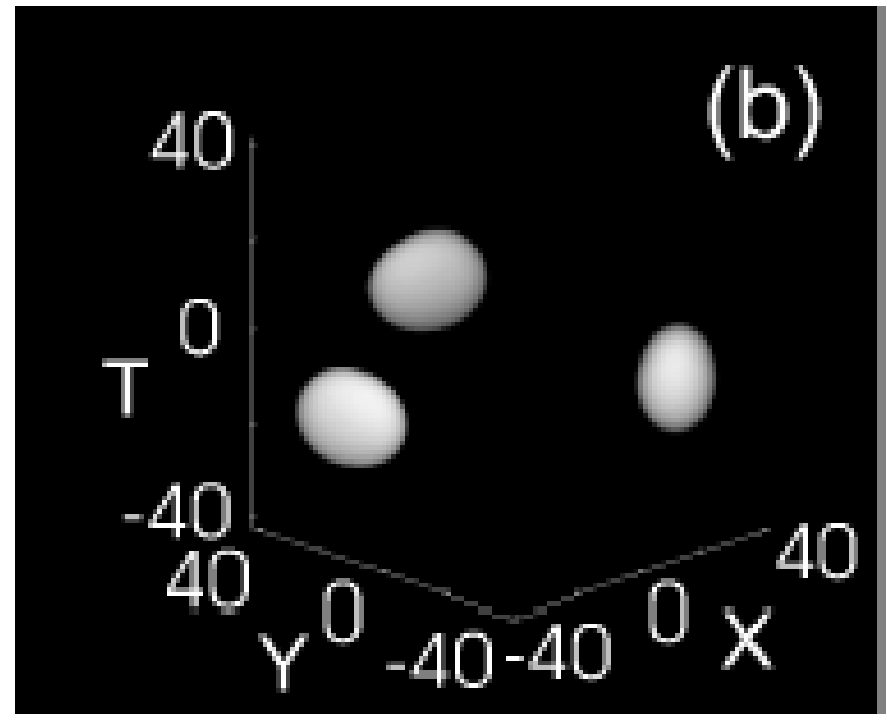
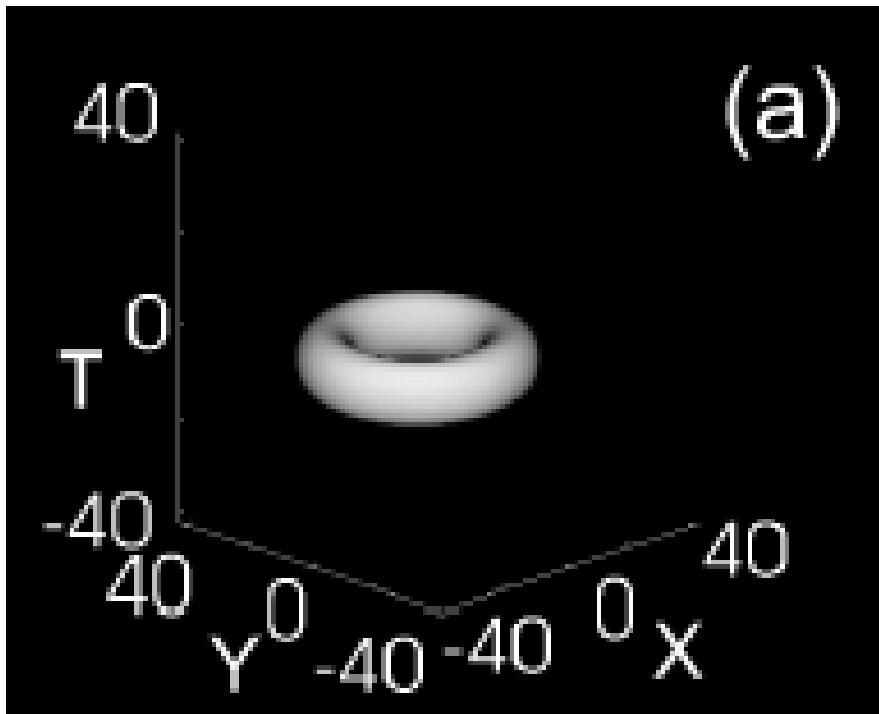
(a)



(b)



An example of the *instability* of the **3D doughnut soliton** with vorticity $S = 2$, viz., *splitting* in *three fragments*:



(4) Novel results: Stable two- and three-dimensional composite solitons in spin-orbit (SO)-coupled self-attractive BEC

(4a) Introduction and objectives

The concept of *emulation* (alias *simulation*) of complex physical effects, known in condensed-matter physics, by much simpler settings available in **BEC** (*matter waves*) and **photonics** (*optical waves*), has drawn a great deal of interest:

P. Hauke, F. M. Cucchietti, L. Tagliacozzo, I. Deutsch, and M. Lewenstein, Rep. Prog. Phys. **75**, 082401 (2012).

An important topic which is enabled by this approach: the emulation of **spin-orbit (SO) interactions** in semiconductors, such as those accounted for by the *Rashba* and *Dresselhaus* Hamiltonians. It is performed by *mapping* the spinor wave function of electrons (*fermions*) into the pseudo-spinor (two-component) *bosonic* wave function of a binary **BEC** gas:

Y. J. Lin, K. Jimenez-Garcia, and I. B. Spielman,
Nature **471**, 83 (2011);

Y. Zhang, L. Mao, and C. Zhang, Phys. Rev. Lett.
108, 035302 (2012);

A review: *H. Zhai*, Rep. Prog. Phys. **78**, 026001
(2015).

(4b) The model

The system of **2D** Gross-Pitaevskii (**GP**) equations for the two-component wave function (ϕ_+, ϕ_-) of the binary **BEC** coupled by the **SO** terms with strength λ and coefficient of the cubic self-attraction $\equiv 1$:

$$i \frac{\partial \phi_+}{\partial t} = -\frac{1}{2} \nabla^2 \phi_+ - |\phi_+|^2 \phi_+ + \lambda \left(\frac{\partial \phi_-}{\partial x} - i \frac{\partial \phi_-}{\partial y} \right),$$
$$i \frac{\partial \phi_-}{\partial t} = -\frac{1}{2} \nabla^2 \phi_- - |\phi_-|^2 \phi_- - \lambda \left(\frac{\partial \phi_+}{\partial x} + i \frac{\partial \phi_+}{\partial y} \right).$$

(4c) References. Basic results are presented here for 2D solitons as per the following papers:

PHYSICAL REVIEW E **89**, 032920 (2014)

Creation of two-dimensional composite solitons in spin-orbit-coupled self-attractive Bose-Einstein condensates in free space

Hidetsugu Sakaguchi and Ben Li

*Department of Applied Science for Electronics and Materials, Interdisciplinary Graduate School of Engineering Sciences,
Kyushu University, Kasuga, Fukuoka 816-8580, Japan*

Boris A. Malomed

Department of Physical Electronics, School of Electrical Engineering, Faculty of Engineering, Tel Aviv University, Tel Aviv 69978, Israel
(Received 12 December 2013; published 26 March 2014)

PHYSICAL REVIEW E **94**, 032202 (2016)

Vortex solitons in two-dimensional spin-orbit coupled Bose-Einstein condensates: Effects of the Rashba-Dresselhaus coupling and Zeeman splitting

Hidetsugu Sakaguchi

*Department of Applied Science for Electronics and Materials, Interdisciplinary Graduate School of Engineering Sciences,
Kyushu University, Kasuga, Fukuoka 816-8580, Japan*

E. Ya. Sherman

*Department of Physical Chemistry, University of the Basque Country UPV-EHU, 48940 Bilbao, Spain
and IKERBASQUE, Basque Foundation for Science, Bilbao, Spain*

Boris A. Malomed

Department of Physical Electronics, School of Electrical Engineering, Faculty of Engineering, Tel Aviv University, Tel Aviv 69978, Israel
(Received 18 February 2016; published 2 September 2016)

(4d) Semi-vortex states

The coupled **GP** equations admit a family of solutions for **semi-vortices**, with vorticities $\mathbf{m}_+ = \mathbf{0}$ in one component, and $\mathbf{m}_- = \mathbf{1}$ in the other. The **exact ansatz** for these solutions, compatible with the underlying equations ($\mu < 0$ is the chemical potential):

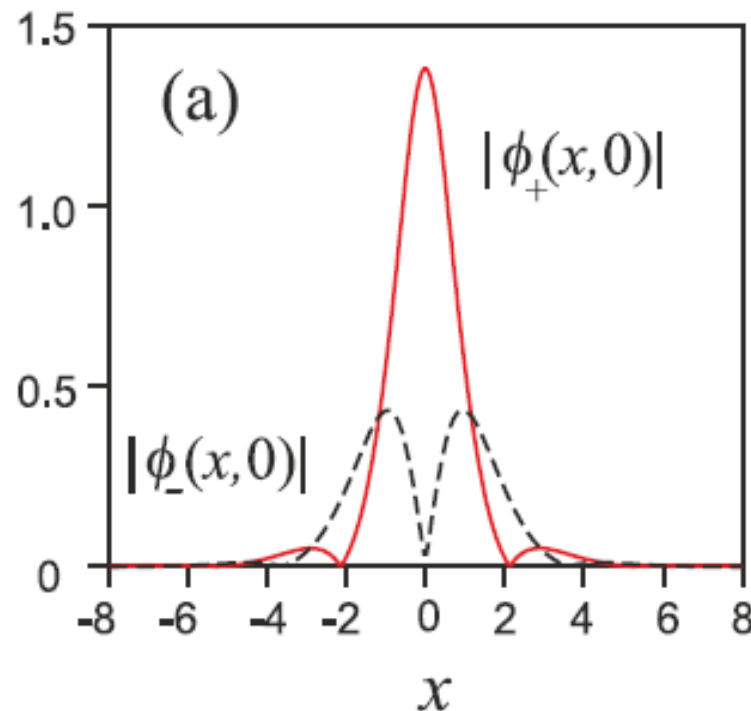
$$\phi_+(x, y, t) = \exp(-i\mu t) f_+(r),$$

$$\phi_-(x, y, t) = \exp(-i\mu t + i\theta) r f_-(r),$$

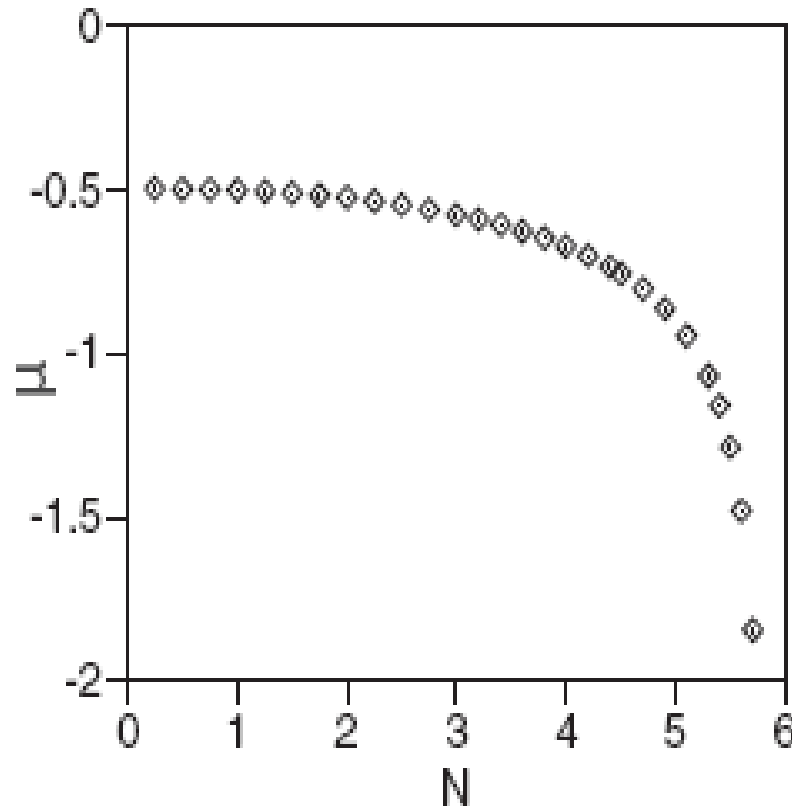
with $f_{\pm}(r)$ taking finite values at $r = 0$

and decaying $\sim \exp(-\sqrt{-2\mu}r)$ at $r \rightarrow \infty$.

A numerically generated cross-section (along $y = 0$) of the two components, $\mathbf{A} \equiv |\boldsymbol{\varphi}_{\pm}|$, for a **stable semi-vortex**:



The numerically found dependence between the total norm of the semi-vortices and their chemical potential demonstrates that **(1)** the norm of the semi-vortex indeed falls **below the threshold value** necessary for the onset of the collapse: $N(\mu) < N_{\text{thr}} \equiv N(\mu \rightarrow -\infty) \approx 5.85$; **(2)** there is **no finite minimum value** of the norm necessary for the existence of the semi-vortex; **(3)** the dependence satisfies the **Vakhitov-Kolokolov** criterion, $d\mu/dN < 0$, which is a necessary condition for the stability:



(4e) Mixed modes

The extended system of the **2D** Gross-Pitaevskii equations for the two-component wave function (ϕ_+, ϕ_-) , which includes both the **linear SO** coupling and **nonlinear cross-interaction** with relative strength $\gamma > 0$:

$$i \frac{\partial \phi_+}{\partial t} = -\frac{1}{2} \nabla^2 \phi_+ - \left(|\phi_+|^2 + \gamma |\phi_-|^2 \right) \phi_+ + \lambda \left(\frac{\partial \phi_-}{\partial x} - i \frac{\partial \phi_-}{\partial y} \right),$$

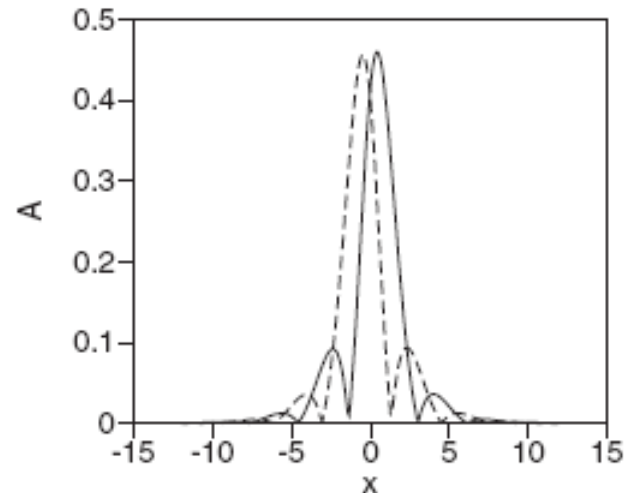
$$i \frac{\partial \phi_-}{\partial t} = -\frac{1}{2} \nabla^2 \phi_- - \left(|\phi_-|^2 + \gamma |\phi_+|^2 \right) \phi_- - \lambda \left(\frac{\partial \phi_+}{\partial x} + i \frac{\partial \phi_+}{\partial y} \right).$$

In the presence of the nonlinear cross-coupling, the semi-vortices remain stable at $\gamma < 1$, and become **unstable** at $\gamma > 1$. In that case, another class of **stable** localized states can be constructed in the form of **mixed modes**, so called because they **mix** fundamental and vortical terms in each component, namely, $m_1 = (0, -1)$ and $m_2 = (0, +1)$, as per the following **ansatz**, which was used as a basis of the **variational approximation** (but it does not produce an exact solution):

$$\phi_+ = A_1 \exp(-\alpha_1 r^2) - A_2 r \exp(-i\theta - \alpha_2 r^2),$$

$$\phi_- = A_1 \exp(-\alpha_1 r^2) + A_2 r \exp(+i\theta - \alpha_2 r^2).$$

A typical example of the cross-section of the mixed mode:

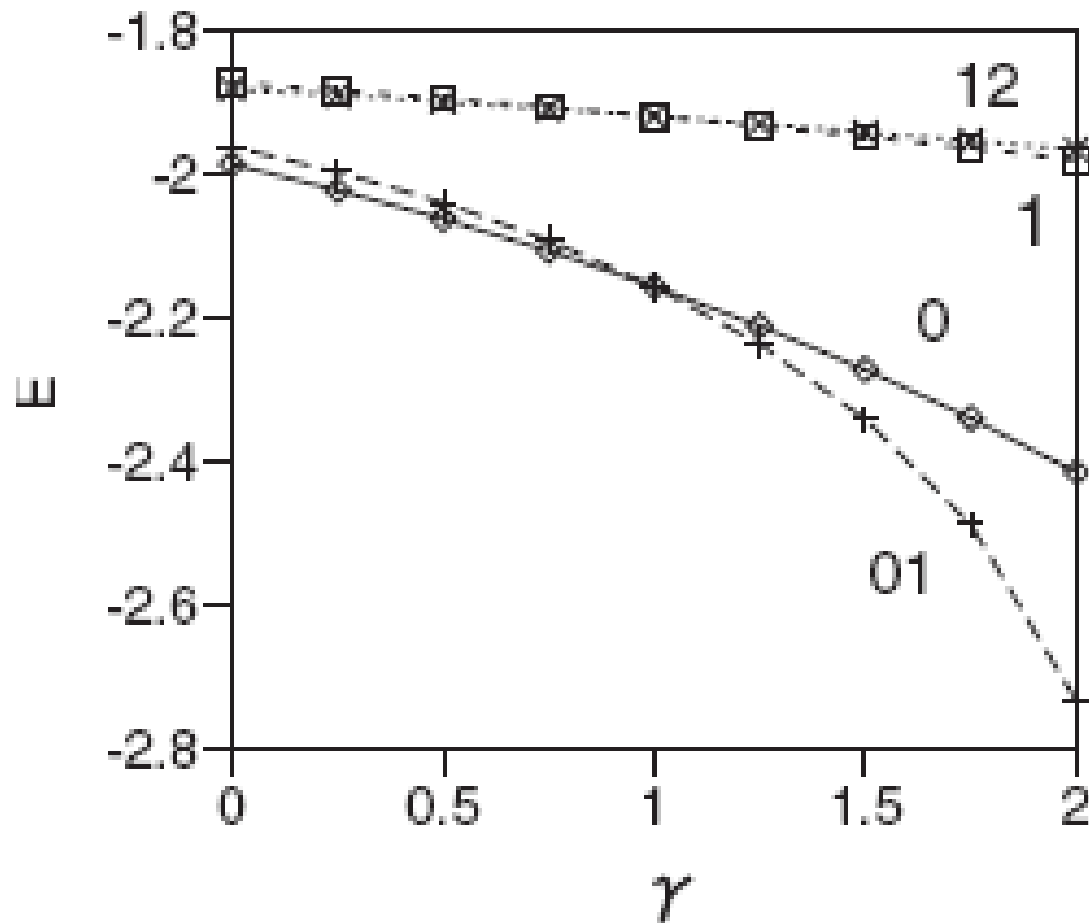


The result is that the **mixed modes** are **unstable** at $\gamma < 1$, and **stable** at $\gamma > 1$, exactly when the semi-vortex is **stable** or **unstable**, respectively.

This **stability switch** between the **semi-vortex** and **mixed mode** with the increase of γ is explained by the fact that the **semi-vortex** and **mixed mode** are **ground states**, which realize the **minimum** of the system's energy, E , precisely at $\gamma < 1$ and $\gamma > 1$, respectively.

$$\begin{aligned}
 E = \iint \bigg\{ & \frac{1}{2}(|\nabla \phi_+|^2 + |\nabla \phi_-|^2) - \frac{1}{2}(|\phi_+|^4 + |\phi_-|^4) \\
 & - \gamma |\phi_+|^2 |\phi_-|^2 + \frac{\lambda}{2} \left[\phi_+^* \left(\frac{\partial \phi_-}{\partial x} - i \frac{\partial \phi_-}{\partial y} \right) \right. \\
 & \left. + \phi_-^* \left(-\frac{\partial \phi_+}{\partial x} - i \frac{\partial \phi_+}{\partial y} \right) \right] + \text{c.c.} \bigg\} dx dy,
 \end{aligned}$$

The dependence of the energy of the **semi-vortex** (“0”) and **mixed mode** (“01”) on γ , for a fixed value of the total norm, $N = 3.7$:



(4f) The stabilization of 3D solitons by the SO coupling

PRL **115**, 253902 (2015)

PHYSICAL REVIEW LETTERS

week ending
18 DECEMBER 2015

Stable Solitons in Three Dimensional Free Space without the Ground State: Self-Trapped Bose-Einstein Condensates with Spin-Orbit Coupling

Yong-Chang Zhang,¹ Zheng-Wei Zhou,^{1,*} Boris A. Malomed,² and Han Pu^{3,4,†}

¹*Key Laboratory of Quantum Information, and Synergetic Innovation Center of Quantum Information and Quantum Physics, University of Science and Technology of China, Hefei, Anhui 230026, China*

²*Department of Physical Electronics, School of Electrical Engineering, Faculty of Engineering, Tel Aviv University, 69978 Tel Aviv, Israel*

³*Department of Physics and Astronomy, and Rice Center for Quantum Materials, Rice University, Houston, Texas 77005, USA*

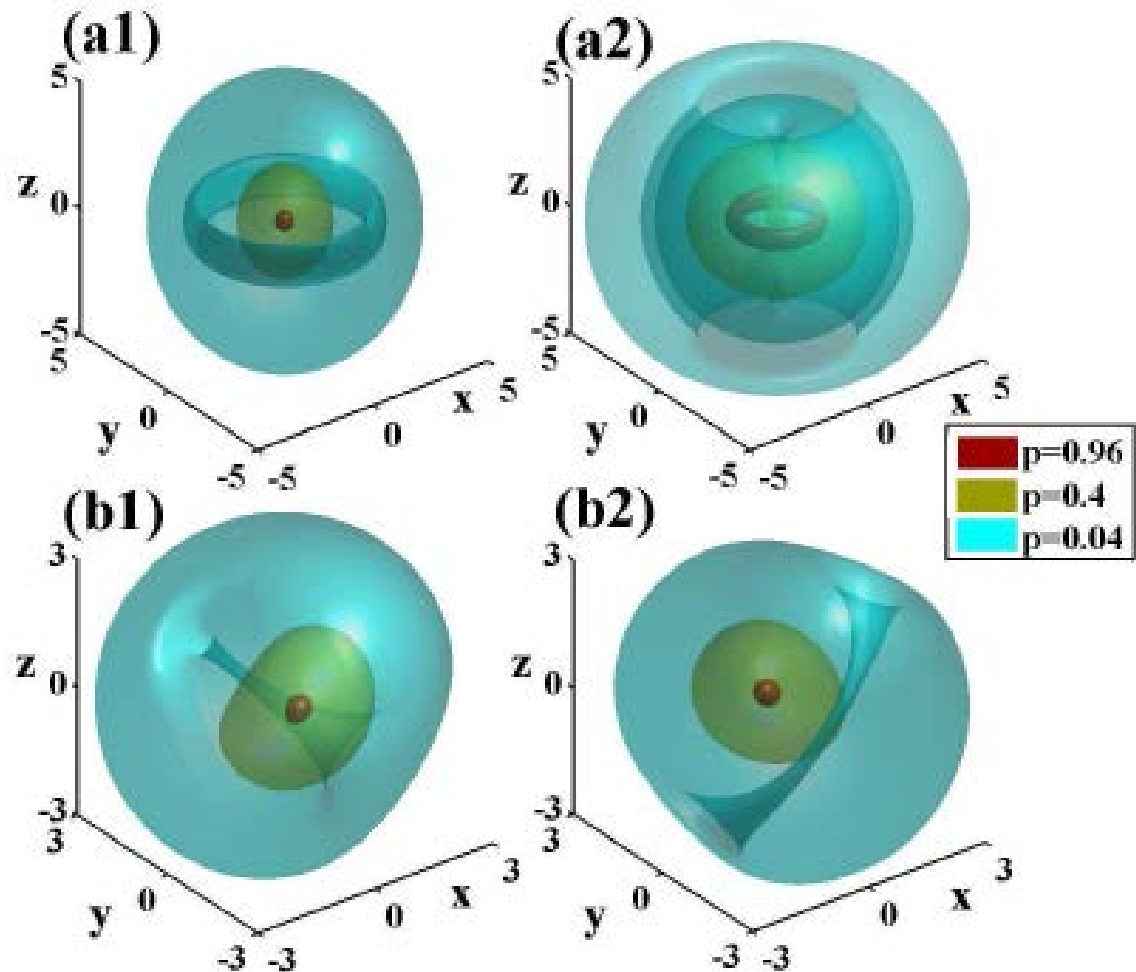
⁴*Center for Cold Atom Physics, Chinese Academy of Sciences, Wuhan 430071, China*

(Received 11 September 2015; published 17 December 2015)

The **3D** model with the **SO** coupling (in particular, of the *Weyl type*) is based on the following system of **GP** equations for the two-component wave function, with the vector of **Pauli matrices**, $\boldsymbol{\sigma} = (\sigma_x, \sigma_y, \sigma_z)$, and $\mathbf{g} \equiv +1$, $\eta > 0$ (the *attractive* nonlinearity, which gives rise to the *supercritical collapse* in **3D**):

$$\left[i \frac{\partial}{\partial t} + \frac{1}{2} \nabla^2 + i \lambda \nabla \cdot \boldsymbol{\sigma} + g \begin{pmatrix} |\psi_+|^2 + \eta |\psi_-|^2 & 0 \\ 0 & |\psi_-|^2 + \eta |\psi_+|^2 \end{pmatrix} \right] \Psi = 0.$$

Examples of (*meta*) *stable* 3D *solitons*, with $N = 8$ [(a), for $\eta = 0.3$ – a semi-vortex; (b), for $\eta = 1.5$ – a *mixed mode*]:

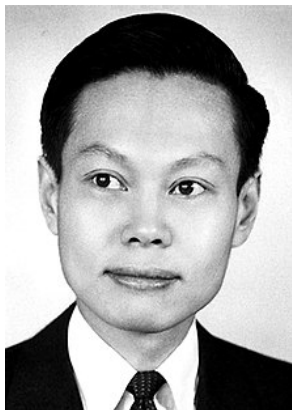


(5) The newest addition: stabilization of 3D and 2D “superfluid droplets” by quantum fluctuations

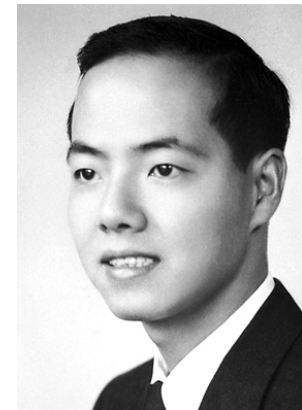
Starting in **2015**, a **stream** of theoretical predictions and *direct experimental observations*, closely related to the topic of multidimensional solitons, has emerged: the creation of **3D** and **2D** soliton-like states in **binary BEC**, named “*quantum droplets*”. They are made **stable** by the **Lee-Huang-Yang (LHY) correction** to the **BEC** dynamics, which represents effects of **quantum fluctuations (Bogoliubov’s modes)** around the **mean-field** states:

T. D. Lee, K. Huang, and C. N. Yang, *Eigenvalues and eigenfunctions of a Bose system of hard spheres and its low-temperature properties*, Phys. Rev. **106**, 1135 (**1957**).

Chen-Ning Yang in 1957
(currently **100** years old)



Tsung-Dao Lee in 1956
(currently **96** years old)



(5a) The idea of the **quantum droplets**, in **3D** and **2D** geometries, was proposed by **Dmitry Petrov** (Paris-Saclay):

PRL 115, 155302 (2015)

PHYSICAL REVIEW LETTERS

week ending
9 OCTOBER 2015

Quantum Mechanical Stabilization of a Collapsing Bose-Bose Mixture

D. S. Petrov

Université Paris-Sud, CNRS, LPTMS, UMR8626, Orsay, F-91405, France

(Received 28 June 2015; published 7 October 2015)

PRL 117, 100401 (2016)

PHYSICAL REVIEW LETTERS

week ending
2 SEPTEMBER 2016

Ultradilute Low-Dimensional Liquids

D. S. Petrov¹ and G. E. Astrakharchik²

¹*LPTMS, CNRS, Univ. Paris Sud, Université Paris-Saclay, 91405 Orsay, France*

²*Departament de Física, Campus Nord B4-B5, Universitat Politècnica de Catalunya, E-08034 Barcelona, Spain*

(Received 24 May 2016; revised manuscript received 28 July 2016; published 1 September 2016)

The starting point is a system of coupled **GP** equations for two components of the mean-field wave function of a **binary BEC**, with **self-repulsion** in each component and **attraction** between the components, which **slightly exceeds** the self-repulsion.

Assuming *equal wave functions* of both components, $\Psi_1 = \Psi_2 \equiv \Psi$, the system of two **GP** equations in **3D** may be reduced to the *single* one:

$$i \frac{\partial \Psi}{\partial t} = -\frac{1}{2} \nabla^2 \Psi - |\Psi|^2 \Psi + \gamma |\Psi|^3 \Psi,$$

where $\gamma > 0$ is an effective strength of the **quartic self-repulsion** term which represents the **LHY correction**.

(5b) Fundamental (zero-vorticity) droplets: solutions and experimental realizations

Stationary solutions for **quantum droplets** with chemical potential $\mu < 0$ are looked for as $\Psi = \exp(-i\mu t) \Phi(r)$, with real functions $\Phi(r)$ satisfying equations

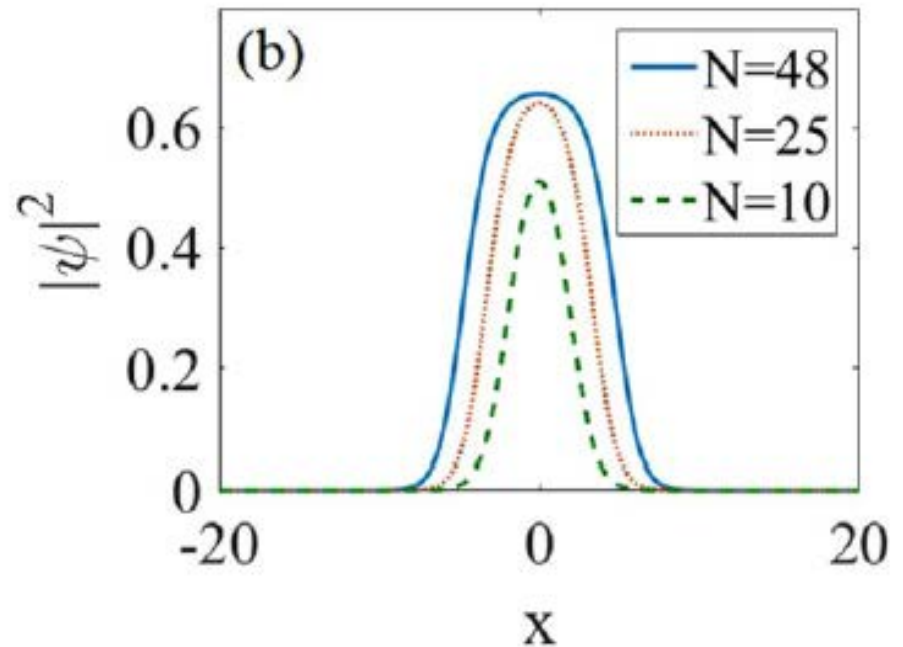
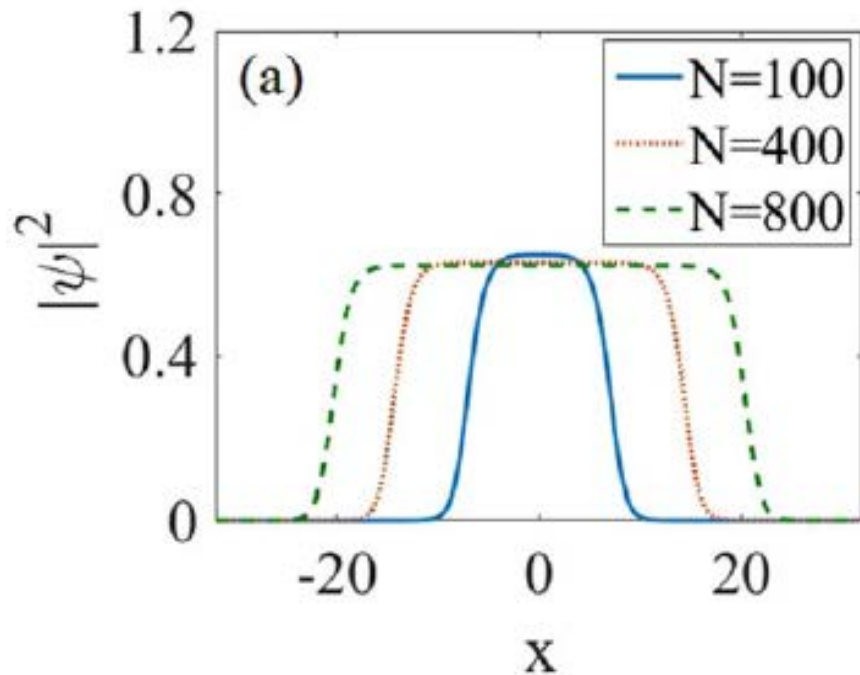
$$\mu \Phi = -\frac{1}{2} \left(\frac{d^2 \Phi}{dr^2} + \frac{2}{r} \frac{d\Phi}{dr} \right) - \Phi^3 + \gamma \Phi^4.$$

Due to the **competition** between the **attractive** and **repulsive** nonlinearities, **density** Φ^2 of the solutions cannot exceed the following **maximum value**:

$$\left(\Phi^2 \right)_{\max} = (25 / 36) \gamma^{-2}.$$

For this reason, the quantum matter described by this equation may be considered as an effectively ***incompressible fluid***, hence the name of “***quantum droplets***”.

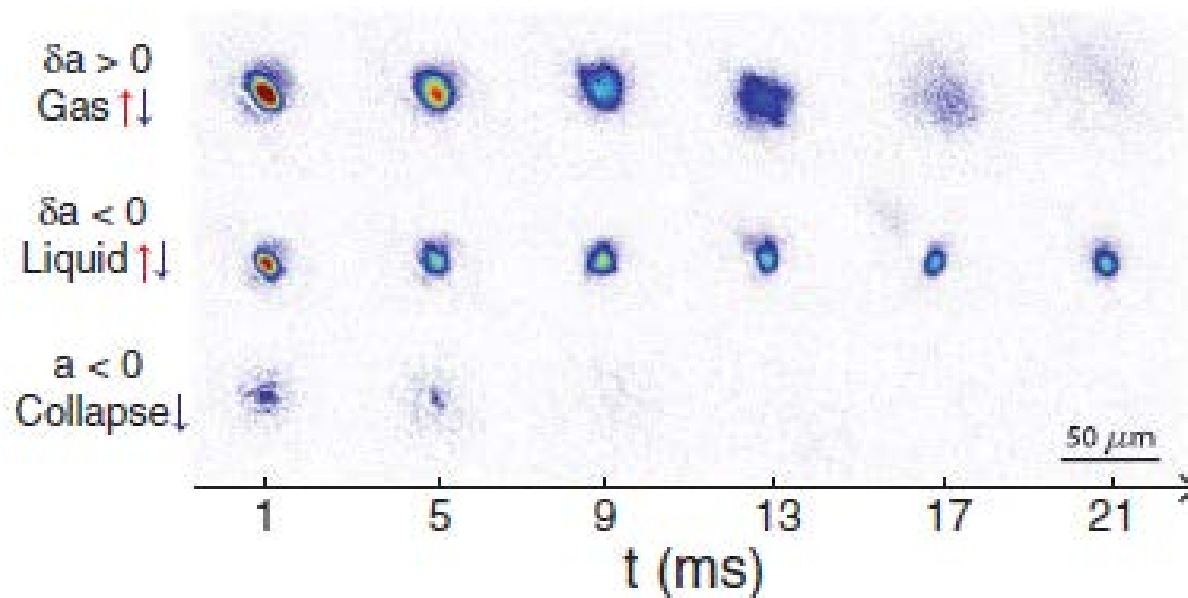
Cross sections of the **quantum-droplet** solutions with **small** (right, **bell-shaped**) and **large** (left, **flat-top** “droplets”) total norms:



Experiment. The creation of **quasi-2D** (*oblate*) droplets, with *aspect ratio* $\sim 10:1$ and lifetime ~ 10 ms, built of $\sim 10,000$ ^{39}K atoms (in two different hyperfine states) per droplet, was reported by the group of **Leticia Tarruell** from ICFO (Barcelona): **Science 359, 301 (2018)**

Quantum liquid droplets in a mixture of Bose-Einstein condensates

C. R. Cabrera,* L. Tanzi,* J. Sanz, B. Naylor, P. Thomas, P. Cheiney, L. Tarruell†



The creation of ***nearly isotropic*** 3D quantum **droplets** (with negligible confinement in any direction) in the same atomic species, ^{39}K , was reported by the group of **Massimo Inguscio** (Florence) :

PHYSICAL REVIEW LETTERS **120**, 235301 (2018)

Self-Bound Quantum Droplets of Atomic Mixtures in Free Space

G. Semeghini,^{1,2,*} G. Ferioli,^{1,2} L. Masi,^{1,2} C. Mazzinghi,¹ L. Wolswijk,¹ F. Minardi,^{2,1,3} M. Modugno,^{4,5}
G. Modugno,^{1,2} M. Inguscio,^{2,1} and M. Fattori^{1,2}

¹*LENS and Dipartimento di Fisica e Astronomia, Università di Firenze, 50019 Sesto Fiorentino, Italy*

²*CNR Istituto Nazionale Ottica, 50019 Sesto Fiorentino, Italy*

³*Dipartimento di Fisica e Astronomia, Università di Bologna, 40127 Bologna, Italy*

⁴*Departamento de Física Teórica e Historia de la Ciencia, Universidad del País Vasco UPV/EHU, 48080 Bilbao, Spain*

⁵*IKERBASQUE, Basque Foundation for Science, 48011 Bilbao, Spain*



(5c) Extension of the theory: quantum droplets with *embedded vorticity*

Vortex-droplet states, with **vorticity** $S = 1, 2, 3, \dots$, were constructed in the **2D** model, looking for solutions (in the polar coordinates, r and θ) as $\Psi = \exp(-i\mu t + iS\theta)U(r)$, with $\mu < 0$.

PHYSICAL REVIEW A **98**, 063602 (2018)

Two-dimensional vortex quantum droplets

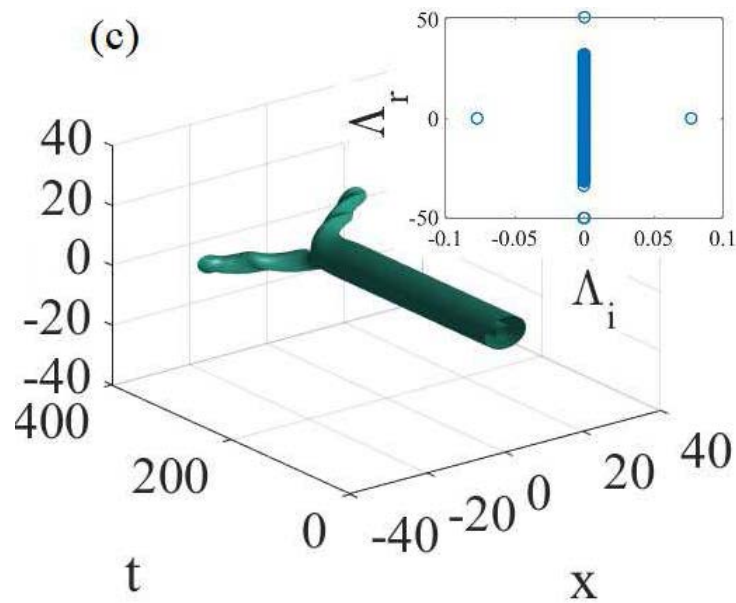
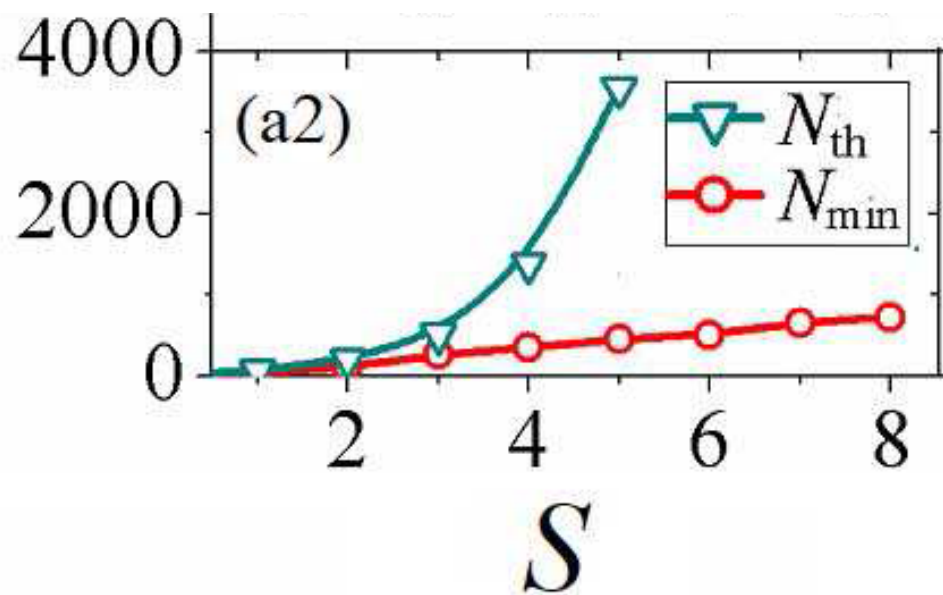
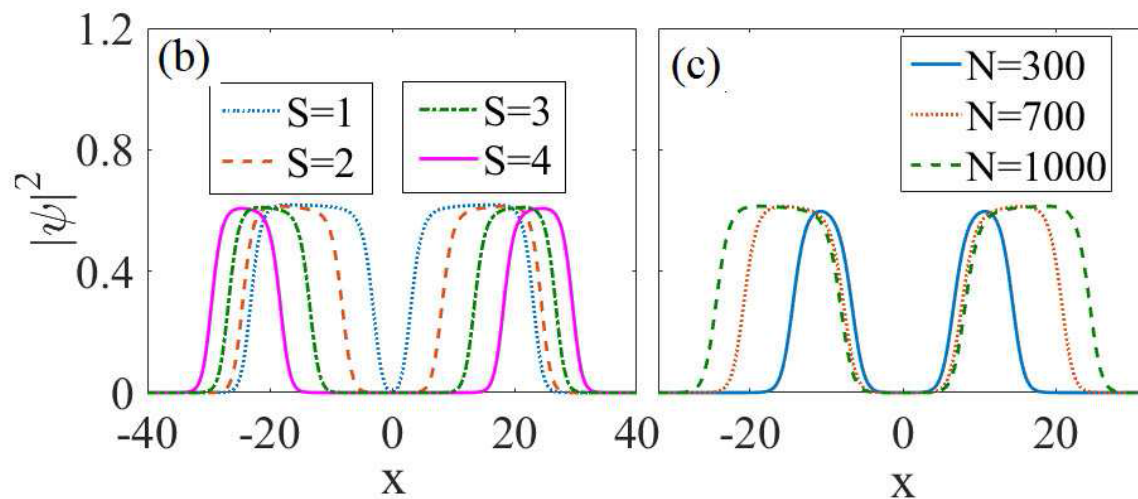
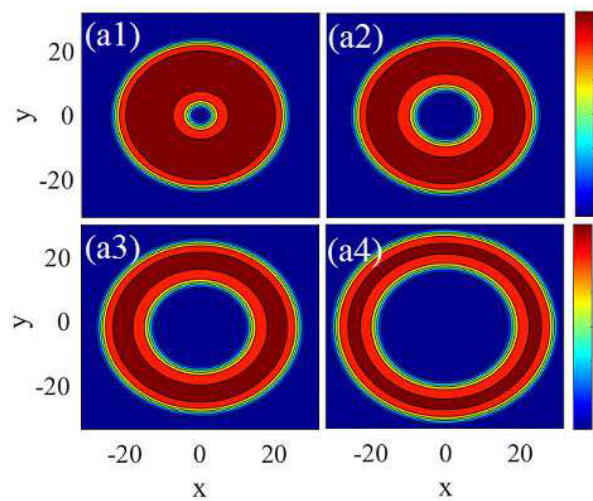
Yongyao Li,¹ Zhaopin Chen,² Zhihuan Luo,³ Chunqing Huang,¹ Haishu Tan,¹ Wei Pang,^{4,*} and Boris A. Malomed^{2,1}

¹*School of Physics and Optoelectronic Engineering, Foshan University, Foshan 528000, China*

²*Department of Physical Electronics, School of Electrical Engineering, Faculty of Engineering, Tel Aviv University, Tel Aviv 69978, Israel*

³*College of Electronic Engineering, South China Agricultural University, Guangzhou 510642, China*

⁴*Department of Experiment Teaching, Guangdong University of Technology, Guangzhou 510006, China*



A more advanced theoretical result: prediction of **stability regions** for **fully 3D** droplets, with embedded vorticities **$S = 1$** and **2** :

PHYSICAL REVIEW A **98**, 013612 (2018)

Three-dimensional droplets of swirling superfluids

Yaroslav V. Kartashov,^{1,2,*} Boris A. Malomed,³ Leticia Tarruell,¹ and Lluís Torner^{1,4}

¹*ICFO-Institut de Ciències Fòniques, The Barcelona Institute of Science and Technology, 08860 Castelldefels, Barcelona, Spain*

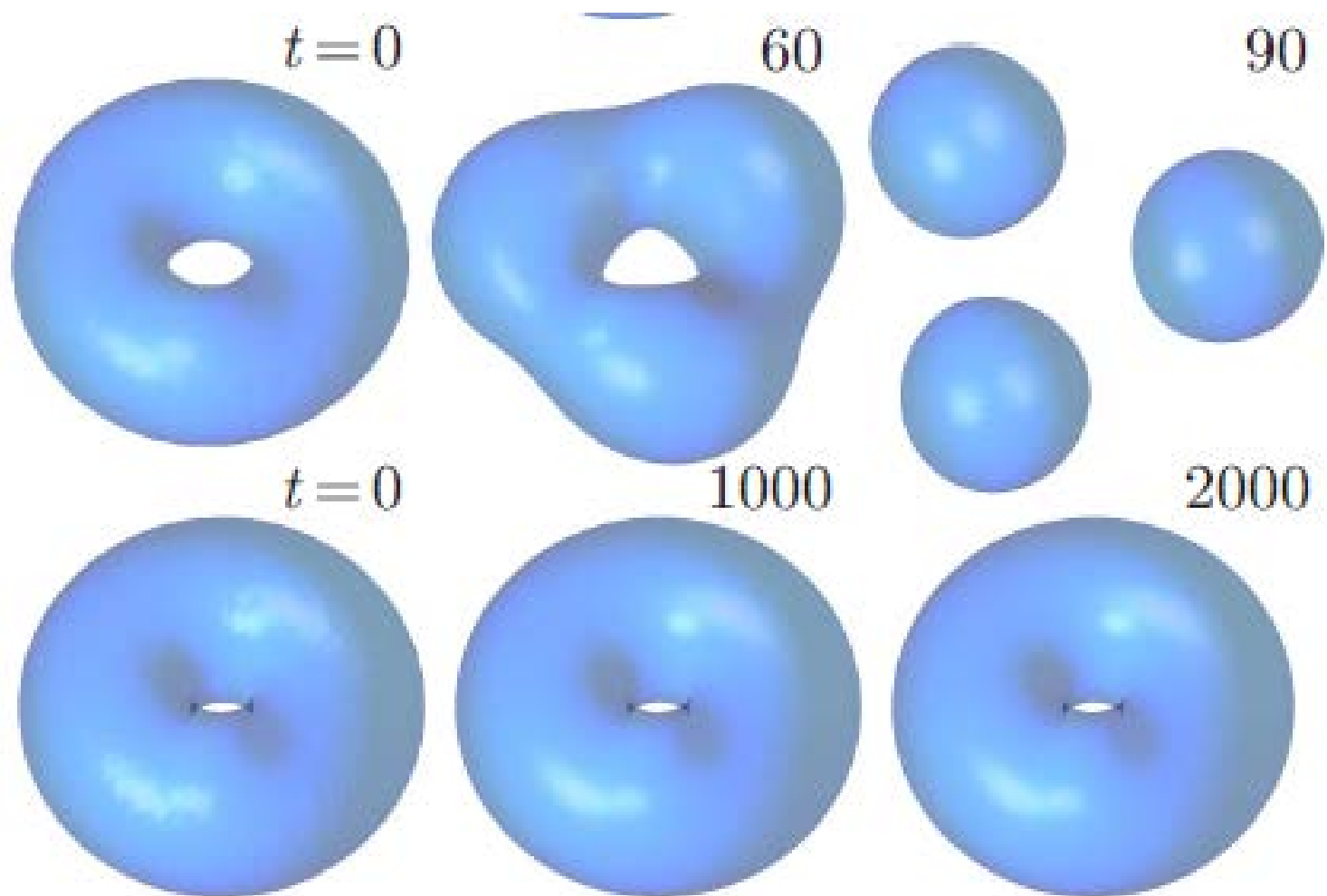
²*Institute of Spectroscopy, Russian Academy of Sciences, Troitsk, Moscow, 108840, Russia*

Department of Physical Electronics, School of Electrical Engineering, Faculty of Engineering, and Center for Light-Matter Interaction, Tel Aviv University, 69978 Tel Aviv, Israel

and ITMO University, St. Petersburg 197101, Russia

⁴*Universitat Politècnica de Catalunya, 08034, Barcelona, Spain*

Examples of *splitting* of an *unstable* 3D *vortex droplet* (the **top row**), and of the evolution of a *stable* one (the **bottom row**). In both cases, $S = 1$.



A relatively recent review on the topic of **quantum droplets** (theory and experiment):



Frontiers of Physics

<https://doi.org/10.1007/s11467-020-1020-2>

Front. Phys.
16(3), 32201 (2021)

REVIEW ARTICLE

A new form of liquid matter: Quantum droplets

Zhi-Huan Luo¹, Wei Pang², Bin Liu³, Yong-Yao Li^{3,†}, Boris A. Malomed^{4,5}

(6) Conclusions

Recent theoretical and experimental studies have led to the prediction and, in some cases, experimental creation of ***stable*** self-trapped modes in the form of **fundamental** and **vortex solitons** in **2D** and **3D** geometries. Especially interesting are the predicted possibilities for the creation of (meta)stable **semi-vortices** by means of the **SO (spin-orbit) coupling** for the binary **BEC** with cubic attractive interactions (something which was previously considered ***absolutely impossible***), as well as the theoretically predicted and ***experimentally realized*** creation of **3D superfluid droplets**, stabilized by **quantum fluctuations**.

Generally, the theoretical studies of multidimensional solitons have advanced much farther than the experimental work. Creation of stable **2D** and **3D** solitons in ***real experiments*** remains a challenging objective.

Thank you for your interest!
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mentioned in it, can be requested from
malomed@tauex.tau.ac.il