

# On nonlocal de Sitter gravity and its cosmological solutions

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(joint work with I. Dimitrijević, B. Dragovich, and J. Stanković)

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- GTR or ETG assumes that Universe is four dimensional homogeneous isotropic simple connected pseudo-Riemannian manifold  $M$  with metric  $(g_{\mu\nu})$  of signature  $(1, 3)$ .
- Generic metric in these spaces is of the form (Friedmann-Robertson-Walker metric (FRW)):

$$ds^2 = -dt^2 + a^2(t) \left( \frac{dr^2}{1 - kr^2} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \right), \quad k \in \{-1, 0, 1\}, \quad (1)$$

where  $a(t)$  is a cosmic scale factor which describes the evolution (in time) of Universe and parameter  $k$  which describes the curvature of the space.

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- GTR is based on Einstein-Hilbert action:

$$S = \int \left( \frac{R - 2\Lambda}{16\pi G c^4} + \mathcal{L}_m \right) \sqrt{-g} d^4x$$

where  $R$  is scalar curvature,  $g = \det(g_{\mu\nu})$  is determinant of metric tensor,  $\Lambda$  is cosmological constant and  $\mathcal{L}_m$  is Lagrangian of matter.

- The variation of the action  $S$  we obtain equations of motion:

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}, \quad c = 1 \quad (2)$$

where  $T_{\mu\nu}$  is the energy momentum tensor,  $g_{\mu\nu}$  is metric tensor,  $R_{\mu\nu}$  is Ricci tensor and  $R$  is scalar curvature.

- The energy momentum tensor for ideal fluid (matter in cosmology) is

$$T = \text{diag}(-\rho g_{00}, g_{11}p, g_{22}p, g_{33}p), \quad (3)$$

where  $\rho$  is energy density and  $p$  is pressure.

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- Einstein equation implies Friedmann equations

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p) + \frac{\Lambda}{3}, \quad \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{k}{a^2} + \frac{\Lambda}{3}.$$

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- Great cosmological observational discoveries of 20th century: high orbital speeds of galaxies in clusters and stars in spiral galaxies, accelerated expansion of the Universe showed that they could not be explained by GTR without additional matter
- Problems related to the Big Bang singularity.

There are two natural approaches:

- Dark matter and energy
- Modification of Einstein theory of gravity, i.e. modification of its Lagrangian  $\mathcal{L}$

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### Dark matter and energy

- Dark matter is responsible for orbital speeds in galaxies, and dark energy is responsible for accelerated expansion of the Universe.
- If Einstein theory of gravity can be applied to the whole Universe then **dark matter and energy** about 5% of ordinary matter, 27% of dark matter and 68% of dark energy.
- It means that 95% of total matter, or energy, represents dark side of the Universe, which nature is unknown.

### Motivation for modification of Einstein theory of gravity

- The validity of General Relativity on cosmological scale is not confirmed.
- Dark matter and dark energy are not yet detected in the laboratory experiments.

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## Different approaches to modification of Einstein theory of gravity

From action

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we have field equations

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where  $T_{\mu\nu}$  is stress-energy tensor,  $g_{\mu\nu}$  is the metric tensor,  $R_{\mu\nu}$  is Ricci tensor and  $R$  is scalar curvature.

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**Theorem 2 (EOM)** The equations of motion for system given by  $S$  are:

$$\tilde{G}_{\mu\nu} = 0, \quad (4)$$

where

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$$\Omega_{\mu\nu} = \sum_{n=1}^{\infty} f_n \sum_{l=0}^{n-1} S_{\mu\nu}(\square^l \mathcal{H}(R), \square^{n-1-l} \mathcal{G}(R)),$$

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- If we take

- $g(R) = \mathcal{H}(R)$  and

- $Q(R)$  be an eigenfunction of the corresponding differential operator operator:  $\square Q(R) = qQ(R)$ , and consequently  $\mathcal{H}'(Q(R)) = \mathcal{H}'(q)Q(R)$ .

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- If we suppose that the manifold  $M$  is endowed with FRW metric, then we have just **two** linearly independent equations: trace and 00-equation.

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for the following cases:

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2.  $\mathcal{H}(R) = R^{-1}$ ,  $\mathcal{G}(R) = R$ ,
3.  $\mathcal{H}(R) = R^p$ ,  $\mathcal{G}(R) = R^q$ ,
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$$S = \int_M \left( \frac{R - 2\Lambda}{16\pi G} + \mathcal{H}(R) \mathcal{F}(\square) \mathcal{G}(R) \right) \sqrt{-g} d^4x,$$

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**3. model:**  $\mathcal{H}(R) = R^p$ ,  $\mathcal{G}(R) = R^q$ ,  $p \geq q$ .

- We considered case with scale factor in the form  $a(t) = a_0 \exp(-\frac{\gamma}{12} t^2)$
- For  $p = q = 1$  there are infinite number of solutions, and constants  $\gamma$  and  $\Lambda$  satisfy  $\gamma = -12\Lambda$ .
- In other cases we proved existence of unique solution, for arbitrary  $\gamma \in \mathbb{R}$ . We explicitly found solutions for  $1 \leq q \leq p \leq 4$ .

**4. model:**  $\mathcal{H}(R) = (R + R_0)^m$ ,  $\mathcal{G}(R) = (R + R_0)^m$ .

- We considered scale factor and ansatz of the form

$$a(t) = At^n \exp(-\frac{\gamma}{12} t^2) \quad \text{and} \quad \square(R + R_0)^m = r(R + R_0)^m.$$

- Using this ansatz we obtained the following five solutions:

- $r = m\gamma$ ,  $n = 0$ ,  $R_0 = \gamma$ ,  $m = \frac{1}{2}$
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**5. model:**  $R = \text{const.}$

- If  $R = R_0 > 0$ , then there exist non-singular solutions for all three values of parameter  $k = 0, \pm 1$ , which are bounced in the cases  $k = 0, 1$ .
- If  $R = R_0 = 0$  then exists Milne's solution  $a(t) = |t + \frac{\pi}{2}|$ .
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- Recently, we have considered the nonlocal gravity model with cosmological constant  $\Lambda$  and without matter, given by

$$(MS) \quad S = \frac{1}{16\pi G} \int_M (R - 2\Lambda + \sqrt{R - 2\Lambda} \mathcal{F}(\square) \sqrt{R - 2\Lambda}) \sqrt{-g} d^4x,$$

where  $\mathcal{F}(\square) = 1 + \sum_{n=1}^{+\infty} f_n \square^n + \sum_{n=1}^{+\infty} l_{-n} \square^{-n}$

- It is a **quasi-linear** since the EOM (5), for  $GG(R) = \sqrt{R - 2\Lambda}$ , is simplified to

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## 1. Cosmological solution in the flat Universe ( $k = 0$ )

### 1.1. Solutions of the form $a(t) = At^\alpha e^{\beta t}$

- There are two solutions:

$$a_1(t) = A t^{\frac{2}{3}} e^{\frac{2}{3}\Lambda t}, \quad \mathcal{F}\left(-\frac{2}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(-\frac{2}{3}\Lambda\right) = 0,$$

$$a_2(t) = A e^{\frac{2}{3}\Lambda t}, \quad \mathcal{F}(-\Lambda) = -1, \quad \mathcal{F}'(-\Lambda) = 0.$$

More solutions of the form  $a(t) = (a_0 + a_1 t) e^{\beta t}$

In this case for  $a_0 \neq 0$ ,  $\beta \neq 0$ , and  $\beta \neq \pm \sqrt{\frac{3}{8}\Lambda}$  we have solutions if

$$\gamma = \frac{2}{3}, \quad \eta = \frac{2}{3}\Lambda, \quad \lambda = \pm \sqrt{\frac{3}{8}\Lambda}.$$

- When  $a_0 \neq 0$ , we have the following two special solutions:

$$a_3(t) = A \cosh^{\frac{2}{3}}\left(\sqrt{\frac{3}{8}\Lambda} t\right), \quad \mathcal{F}\left(\frac{3}{8}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{3}{8}\Lambda\right) = 0,$$

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## 1. Cosmological solution in the flat Universe ( $k = 0$ )

### 1.1. Solutions of the form $a(t) = A t^n e^{\gamma t^2}$

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$$a_1(t) = A t^{\frac{3}{2}} e^{\frac{\Lambda}{12} t^2}, \quad \mathcal{F}\left(-\frac{3}{7}\Lambda\right) = -1, \quad \mathcal{F}'\left(-\frac{3}{7}\Lambda\right) = 0,$$

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### 1.2. New solutions of the form $a(t) = (\alpha e^{\lambda t} + \beta e^{-\lambda t})^\gamma$

- In this case for  $\alpha\beta \neq 0$ ,  $R \neq 2\Lambda$  and  $q \neq 0$  we have solutions if

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### 1.2. New solutions of the form $a(t) = (\alpha e^{\lambda t} + \beta e^{-\lambda t})^\gamma$

- In this case for  $\alpha\beta \neq 0$ ,  $R \neq 2\Lambda$  and  $q \neq 0$  we have solutions if

$$\gamma = \frac{2}{3}, \quad q = \frac{3}{8}\Lambda, \quad \lambda = \pm\sqrt{\frac{3}{8}\Lambda}.$$

- When  $\alpha\beta \neq 0$ , we have the following two special solutions:

$$a_3(t) = A \cosh^{\frac{2}{3}}\left(\sqrt{\frac{3}{8}\Lambda} t\right), \quad \mathcal{F}\left(\frac{3}{8}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{3}{8}\Lambda\right) = 0,$$

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## 1. Cosmological solution in the flat Universe ( $k = 0$ )

### 1.1. Solutions of the form $a(t) = A t^n e^{\gamma t^2}$

- There are two solutions:

$$a_1(t) = A t^{\frac{2}{3}} e^{\frac{\Lambda}{14} t^2}, \quad \mathcal{F}\left(-\frac{3}{7}\Lambda\right) = -1, \quad \mathcal{F}'\left(-\frac{3}{7}\Lambda\right) = 0,$$

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## 1. Cosmological solution in the flat Universe ( $k = 0$ )

### 1.3. New solutions of the form $a(t) = (\alpha \sin M + \beta \cos M)^2$

- For  $\alpha \neq 0$  and  $\beta \neq 0$  there are only possibility for  $\gamma, \gamma' = \frac{3}{8}$ . Taking  $A = a_0^2$  and  $\Lambda = -\frac{3}{8}A$ , we have the following two solutions:

$$a(t) = A \left( 1 + \sin^2 \left( \sqrt{-\frac{3}{8}\Lambda} t \right) \right), \quad \mathcal{F}\left(\frac{3}{8}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{3}{8}\Lambda\right) = 0$$

$$a(t) = A \left( 1 - \sin^2 \left( \sqrt{-\frac{3}{8}\Lambda} t \right) \right), \quad \mathcal{F}\left(\frac{3}{8}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{3}{8}\Lambda\right) = 0$$

- For  $\alpha = 0$  or  $\beta = 0$ , we have also two cosmological solutions with  $\gamma = \frac{3}{8}$ :

$$a(t) = A \sin^2 \left( \sqrt{-\frac{3}{8}\Lambda} t \right), \quad \mathcal{F}\left(\frac{3}{8}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{3}{8}\Lambda\right) = 0$$

$$a(t) = A \cos^2 \left( \sqrt{-\frac{3}{8}\Lambda} t \right), \quad \mathcal{F}\left(\frac{3}{8}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{3}{8}\Lambda\right) = 0$$

## 1. Cosmological solution in the flat Universe ( $k = 0$ )

### 1.3. New solutions of the form $a(t) = (\alpha \sin \lambda t + \beta \cos \lambda t)^\gamma$

- For  $\alpha \neq 0$  and  $\beta \neq 0$  there are only possibility for  $\gamma$ ,  $\gamma = \frac{2}{3}$ . Taking  $\beta = \pm\alpha$ , and  $A = \alpha^{\frac{3}{2}}$ , we have the following two solutions:

$$a_5(t) = A \left( 1 + \sin \left( 2\sqrt{-\frac{3}{8}\Lambda} t \right) \right)^{\frac{1}{3}}, \quad \mathcal{F}\left(\frac{3}{8}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{3}{8}\Lambda\right) = 0,$$

$$a_6(t) = A \left( 1 - \sin \left( 2\sqrt{-\frac{3}{8}\Lambda} t \right) \right)^{\frac{1}{3}}, \quad \mathcal{F}\left(\frac{3}{8}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{3}{8}\Lambda\right) = 0.$$

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$$a_7(t) = A \sin^{\frac{2}{3}} \left( \sqrt{-\frac{3}{8}\Lambda} t \right), \quad \mathcal{F}\left(\frac{3}{8}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{3}{8}\Lambda\right) = 0,$$

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## 2. Cosmological solution in the open and closed Universe ( $k = \pm 1$ )

### 2.1. Solutions of the form $a(t) = A e^{\pm \sqrt{\Lambda} t}$ , ( $k = \pm 1$ )

- For  $\alpha \neq 0, \beta = 0$  or  $\alpha = 0, \beta \neq 0$  we have the following solution:

$$a_0(t) = A e^{\pm \sqrt{\Lambda} t}, \quad k = \pm 1, \quad \mathcal{F}\left(\frac{1}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{1}{3}\Lambda\right) = 0, \quad A > 0$$

- Here solutions of the form  $a(t) = (A \cosh^2 \pm B \sinh^2) e^{\pm \sqrt{\Lambda} t}$ , ( $k = \pm 1$ )

- For  $\alpha \neq 0, \beta \neq 0, \beta \neq 2\Lambda, \alpha \neq 0$  there are two following cosmological solutions:

$$a_0(t) = A \cosh^2\left(\sqrt{\frac{2}{3}}\Lambda t\right), \quad k = \pm 1, \quad \mathcal{F}\left(\frac{1}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{1}{3}\Lambda\right) = 0,$$

$$a_1(t) = A \sinh^2\left(\sqrt{\frac{2}{3}}\Lambda t\right), \quad k = \pm 1, \quad \mathcal{F}\left(\frac{1}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{1}{3}\Lambda\right) = 0.$$

## 2. Cosmological solution in the open and closed Universe ( $k = \pm 1$ )

2.1. Solutions of the form  $a(t) = A e^{\pm\sqrt{\frac{\Lambda}{3}}t}$ , ( $k = \pm 1$ )

- For  $\alpha \neq 0, \beta = 0$  or  $\alpha = 0, \beta \neq 0$  we have the following solution:

$$a_0(t) = A e^{\pm\sqrt{\frac{\Lambda}{3}}t}, \quad k = \pm 1, \quad \mathcal{F}\left(\frac{1}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{1}{3}\Lambda\right) = 0, \quad \Lambda > 0.$$

2.2. New solutions of the form  $a(t) = (\alpha e^{\lambda t} + \beta e^{-\lambda t})^\gamma$ , ( $k = \pm 1$ )

- For  $\alpha \neq 0, \beta \neq 0, R \neq 2\Lambda, q \neq 0$  there are two following cosmological solutions:

$$a_{10}(t) = A \cosh^{\frac{1}{2}}\left(\sqrt{\frac{2}{3}\Lambda}t\right), \quad k = \pm 1, \quad \mathcal{F}\left(\frac{1}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{1}{3}\Lambda\right) = 0,$$

$$a_{11}(t) = A \sinh^{\frac{1}{2}}\left(\sqrt{\frac{2}{3}\Lambda}t\right), \quad k = \pm 1, \quad \mathcal{F}\left(\frac{1}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{1}{3}\Lambda\right) = 0.$$

## 2. Cosmological solution in the open and closed Universe ( $k = \pm 1$ )

2.1. Solutions of the form  $a(t) = A e^{\pm\sqrt{\frac{\Lambda}{6}}t}$ , ( $k = \pm 1$ )

- For  $\alpha \neq 0, \beta = 0$  or  $\alpha = 0, \beta \neq 0$  we have the following solution:

$$a_9(t) = A e^{\pm\sqrt{\frac{\Lambda}{6}}t}, \quad k = \pm 1, \quad \mathcal{F}\left(\frac{1}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{1}{3}\Lambda\right) = 0, \quad \Lambda > 0.$$

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## 2. Cosmological solution in the open and closed Universe ( $k = \pm 1$ )

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$$a_{10}(t) = A \cosh^{\frac{1}{2}}\left(\sqrt{\frac{2}{3}\Lambda}t\right), \quad k = \pm 1, \quad \mathcal{F}\left(\frac{1}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{1}{3}\Lambda\right) = 0,$$

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- For  $\alpha \neq 0, \beta \neq 0, R \neq 2\Lambda, q \neq 0$  there are two following cosmological solutions:

$$a_{10}(t) = A \cosh^{\frac{1}{2}}\left(\sqrt{\frac{2}{3}\Lambda}t\right), \quad k = \pm 1, \quad \mathcal{F}\left(\frac{1}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{1}{3}\Lambda\right) = 0,$$

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## 2. Cosmological solution in the open and closed Universe ( $k = \pm 1$ )

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- For  $\alpha \neq 0, \beta \neq 0, R \neq 2\Lambda, q \neq 0$  there are two following cosmological solutions:

$$a_{10}(t) = A \cosh^{\frac{1}{2}}\left(\sqrt{\frac{2}{3}\Lambda}t\right), \quad k = \pm 1, \quad \mathcal{F}\left(\frac{1}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{1}{3}\Lambda\right) = 0,$$

$$a_{11}(t) = A \sinh^{\frac{1}{2}}\left(\sqrt{\frac{2}{3}\Lambda}t\right), \quad k = \pm 1, \quad \mathcal{F}\left(\frac{1}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{1}{3}\Lambda\right) = 0.$$

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2.2. New solutions of the form  $a(t) = (\alpha e^{\lambda t} + \beta e^{-\lambda t})^\gamma$ , ( $k = \pm 1$ )

- For  $\alpha \neq 0, \beta \neq 0, R \neq 2\Lambda, q \neq 0$  there are two following cosmological solutions:

$$a_{10}(t) = A \cosh^{\frac{1}{2}}\left(\sqrt{\frac{2}{3}\Lambda}t\right), \quad k = \pm 1, \quad \mathcal{F}\left(\frac{1}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{1}{3}\Lambda\right) = 0,$$

$$a_{11}(t) = A \sinh^{\frac{1}{2}}\left(\sqrt{\frac{2}{3}\Lambda}t\right), \quad k = \pm 1, \quad \mathcal{F}\left(\frac{1}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{1}{3}\Lambda\right) = 0.$$

## 2. Cosmological solution in the open and closed Universe ( $k = \pm 1$ )

2.1. Solutions of the form  $a(t) = A e^{\pm\sqrt{\frac{\Lambda}{6}}t}$ , ( $k = \pm 1$ )

- For  $\alpha \neq 0, \beta = 0$  or  $\alpha = 0, \beta \neq 0$  we have the following solution:

$$a_9(t) = A e^{\pm\sqrt{\frac{\Lambda}{6}}t}, \quad k = \pm 1, \quad \mathcal{F}\left(\frac{1}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{1}{3}\Lambda\right) = 0, \quad \Lambda > 0.$$

2.2. New solutions of the form  $a(t) = (\alpha e^{\lambda t} + \beta e^{-\lambda t})^\gamma$ , ( $k = \pm 1$ )

- For  $\alpha \neq 0, \beta \neq 0, R \neq 2\Lambda, q \neq 0$  there are two following cosmological solutions:

$$a_{10}(t) = A \cosh^{\frac{1}{2}}\left(\sqrt{\frac{2}{3}\Lambda} t\right), \quad k = \pm 1, \quad \mathcal{F}\left(\frac{1}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{1}{3}\Lambda\right) = 0,$$

$$a_{11}(t) = A \sinh^{\frac{1}{2}}\left(\sqrt{\frac{2}{3}\Lambda} t\right), \quad k = \pm 1, \quad \mathcal{F}\left(\frac{1}{3}\Lambda\right) = -1, \quad \mathcal{F}'\left(\frac{1}{3}\Lambda\right) = 0.$$



- 1. Cosmological solution for  $a_1(t) = At^{\frac{3}{2}} e^{\frac{\Lambda}{14}t^2}$ ,  $k=0$
- The corresponding  $\rho$ , acceleration and the scalar  $R$  curvature are:

$$H_1(t) = \frac{\dot{a}_1}{a_1} = \frac{2}{3} \frac{1}{t} + \frac{1}{7} \Lambda t,$$

$$\ddot{a}_1(t) = \left( -\frac{2}{9} \frac{1}{t^2} + \frac{1}{3} \Lambda + \frac{1}{49} \Lambda^2 t^2 \right) a_1(t),$$

$$R_1(t) = \frac{4}{3} \frac{1}{t^2} + \frac{22}{7} \Lambda + \frac{12}{49} \Lambda^2 t^2,$$

- Friedman equations gives

$$\bar{\rho}(t) = \frac{2t^{-2} + \frac{9}{98} \Lambda^2 t^2 - \frac{9}{14} \Lambda}{12\pi G}, \quad \bar{p}(t) = -\frac{\Lambda}{56\pi G} \left( \frac{3}{7} \Lambda t^2 - 1 \right), \quad (7)$$

where  $\bar{\rho}$  and  $\bar{p}$  are analogs of the energy density and pressure of the dark side of the universe, respectively. The corresponding equation of state is  $\bar{p}(t) = \bar{w}(t) \bar{\rho}(t)$ .

- 1. Cosmological solution for  $a_1(t) = A t^{\frac{2}{3}} e^{\frac{\Lambda}{14} t^2}$ ,  $k = 0$
- The corresponding Hubble parameter, acceleration and the scalar 2 curvature are:

$$H_1(t) = \frac{\dot{a}_1}{a_1} = \frac{2}{3} \frac{1}{t} + \frac{1}{7} \Lambda t,$$

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- Equation (7) implies that  $\tilde{w}(t) \rightarrow -1$  when  $t \rightarrow \infty$ , what corresponds to an analog of  $\Lambda$  dark energy dominance in the standard cosmological model.
- It means that this nonlocal gravity model with cosmological solution  $a(t) = A t^{\frac{2}{3}} e^{\frac{\Lambda}{3} t^2}$  describes some effects usually attributed to the dark matter and dark energy.
- This solution is invariant under transformation  $t \rightarrow -t$  and singular at cosmic time  $t = 0$ .
- Let us recall, the second Friedman equation

$$H^2 = \frac{\dot{a}^2}{a^2} = \frac{8\pi G}{3} \rho - \frac{k}{a^2} + \frac{\Lambda}{3}, \quad (8)$$

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- Then we can rewrite the previous equation as,

$$\begin{aligned} H^2 &= \frac{\dot{a}^2}{a^2} = \frac{8\pi G}{3}\rho_r + \frac{8\pi G}{3}\rho_m - \frac{k}{a^2} + \frac{\Lambda}{3} \\ &= \frac{8C_r\pi G}{a^4} + \frac{8C_m\pi G}{a^3} - \frac{k}{a^2} + \frac{\Lambda}{3} \end{aligned}$$

- or,

$$\frac{H^2}{H_0^2} = \frac{\Omega_r}{a^4} + \frac{\Omega_m}{a^3} + \frac{\Omega_k}{a^2} + \Omega_\Lambda$$

- Observational data obtained by Planck-2018 for the  $\Lambda$ CDM model:

$t_0 = (13.801 \pm 0.024) \times 10^9 \text{yr}$  – age of the universe,

$H(t_0) = (67.40 \pm 0.50) \text{ km/s/Mpc}$  – Hubble parameter,

$\Omega_m = 0.315 \pm 0.007$ – matter density parameter,

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$$H_1(t) = \frac{\dot{a}_1}{a_1} = \frac{2}{3} \frac{1}{t} + \frac{1}{7} \Lambda t,$$

taking  $H_1(t_0) = H(t_0)$  we calculate  $\Lambda_1 = 1.05 \times 10^{-35} \text{s}^{-2}$  that differs from  $\Lambda = 3H^2(t_0) \Omega_\Lambda = 0.98 \times 10^{-35} \text{s}^{-2}$  (by  $\Lambda$ CDM model).

- We also computed

$$\ddot{a}_1(t_0)/a_1(t_0) = 2.7 \times 10^{-36} \text{s}^{-2}$$

$$R(t_0) = 4.5 \times 10^{-35} \text{s}^{-2} \quad \text{and consequently}$$

$$R(t_0) - 2\Lambda = 2.4 \times 10^{-35} \text{s}^{-2}.$$

- Replacing solution  $a_1(t)$  with  $k = 0$ , Friedman equations give

$$\bar{\rho}_1(t) = \frac{3}{8\pi G} \left( H_1^2(t) - \frac{\Lambda_1}{3} \right) = \frac{3}{8\pi G} \left( \frac{4}{9} t^{-2} - \frac{1}{7} \Lambda_1 + \frac{1}{49} \Lambda_1^2 t^2 \right),$$

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- For  $t = t_0$ , from previous formula, and from  $\Lambda$ CDM model we have

$$\bar{\rho}_1(t_0) = 2.26 \times 10^{-30} \frac{g}{cm^3},$$

$$\rho(t_0) = \frac{3}{8\pi G} \left( H_0^2 - \frac{\Lambda}{3} \right) = 2.68 \times 10^{-30} \frac{g}{cm^3}.$$

- Then, for vacuum energy density of background solution  $a_1(t)$  and  $\Lambda$ CDM model, we have

$$\rho(t_0) - \bar{\rho}_1(t_0) = \frac{\Lambda_1 - \Lambda}{8\pi G} = \rho_{\Lambda_1} - \rho_{\Lambda} = 0.42 \times 10^{-30} \frac{g}{cm^3},$$

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- Effective pressure. At the beginning,  $\bar{p}_1(0) = \frac{\Lambda_1}{56\pi G} > 0$ , then decreases and equals zero at  $t = \sqrt{\frac{7}{3\Lambda_1}} = 4,71 \times 10^{17} \text{ s} = 14,917 \times 10^9 \text{ yr}$ .
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**THANK YOU FOR  
YOUR ATTENTION !!!**

Non-trivial Christoffel symbols of Friedman – Robertson – Walker metric

$$\Gamma_{01}^1 = \frac{\dot{a}}{a}$$

$$\Gamma_{02}^2 = \frac{\dot{a}}{a}$$

$$\Gamma_{03}^3 = \frac{\dot{a}}{a}$$

$$\Gamma_{11}^0 = \frac{a \dot{a}}{1 - k r^2}$$

$$\Gamma_{11}^1 = \frac{k r}{1 - k r^2}$$

$$\Gamma_{12}^2 = \frac{1}{r}$$

$$\Gamma_{13}^3 = \frac{1}{r}$$

$$\Gamma_{22}^0 = r^2 a \dot{a}$$

$$\Gamma_{22}^1 = r(k r^2 - 1)$$

$$\Gamma_{23}^3 = \cot \theta$$

$$\Gamma_{33}^0 = r^2 a \dot{a} \sin^2 \theta$$

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Non-trivial components of curvature tensor

$$\begin{aligned}
 R_{0110} &= \frac{a \ddot{a}}{1 - k r^2} & R_{1221} &= -\frac{r^2 a^2 (\dot{a}^2 + k)}{1 - k r^2} \\
 R_{0220} &= r^2 a \ddot{a} & R_{1331} &= -\frac{r^2 a^2 \sin^2 \theta (\dot{a}^2 + k)}{1 - k r^2} \\
 R_{0330} &= r^2 a \ddot{a} \sin^2 \theta & R_{2332} &= -r^4 a^2 \sin^2 \theta (\dot{a}^2 + k)
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Ricci tensor

$$R_{\mu\nu} = \begin{pmatrix} -\frac{3\ddot{a}}{a} & 0 & 0 & 0 \\ 0 & u g_{11} & 0 & 0 \\ 0 & 0 & u g_{22} & 0 \\ 0 & 0 & 0 & u g_{33} \end{pmatrix}, \quad u = \frac{a \ddot{a} + 2(\dot{a}^2 + k)}{a^2}$$

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Scalar curvature

$$R = \frac{6(a\ddot{a} + \dot{a}^2 + k)}{a^2}$$

Einstein tensor

$$G_{\mu\nu} = \begin{pmatrix} \frac{3(\dot{a}^2 + k)}{a^2} & 0 & 0 & 0 \\ 0 & -v g_{11} & 0 & 0 \\ 0 & 0 & -v g_{22} & 0 \\ 0 & 0 & 0 & -v g_{33} \end{pmatrix}, \quad v = \frac{2a\ddot{a} + \dot{a}^2 + k}{a^2}$$

► FRW metric

► EOM

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$$R = \frac{6(a\ddot{a} + \dot{a}^2 + k)}{a^2}$$

Einstein tensor

$$G_{\mu\nu} = \begin{pmatrix} \frac{3(\dot{a}^2 + k)}{a^2} & 0 & 0 & 0 \\ 0 & -v g_{11} & 0 & 0 \\ 0 & 0 & -v g_{22} & 0 \\ 0 & 0 & 0 & -v g_{33} \end{pmatrix}, \quad v = \frac{2a\ddot{a} + \dot{a}^2 + k}{a^2}$$

► FRW metric

► EOM

► EOM 2

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▶ FRW metric

▶ EOM

▶ EOM-2